

ON THE DUALITY OF TELEPORTATION AND DENSE CODING

arXiv:2302.14798

Felix Leditzky

University of Illinois Urbana-Champaign

July 27, 2023



Eric Chitambar



TQC 2023, Aveiro, Portugal

Quantum teleportation

PHYSICAL REVIEW LETTERS

VOLUME 70

29 MARCH 1993

NUMBER 13

Teleporting an Unknown Quantum State via Dual Classical and Einstein-Podolsky-Rosen Channels

Charles H. Bennett,⁽¹⁾ Gilles Brassard,⁽²⁾ Claude Crépeau,^{(2),(3)}
Richard Jozsa,⁽²⁾ Asher Peres,⁽⁴⁾ and William K. Wootters⁽⁵⁾

⁽¹⁾ IBM Research Division, T.J. Watson Research Center, Yorktown Heights, New York 10598

⁽²⁾ Département IRO, Université de Montréal, C.P. 6128, Succursale "A", Montréal, Québec, Canada H3C 3J7

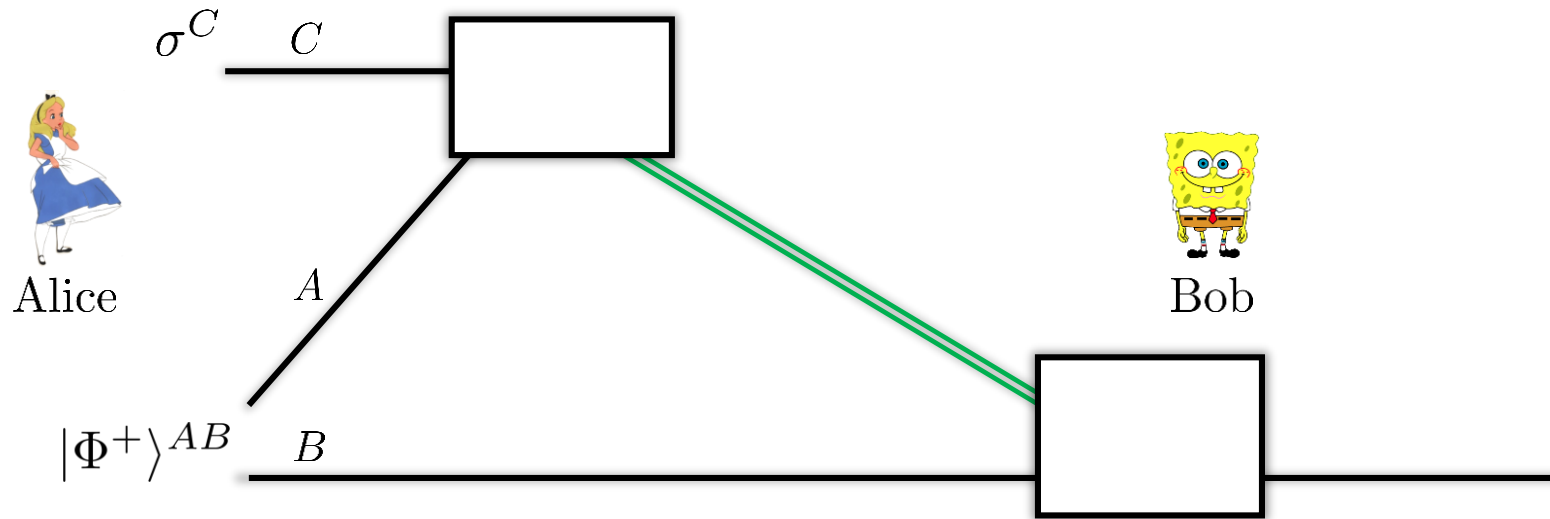
⁽³⁾ Laboratoire d'Informatique de l'École Normale Supérieure, 45 rue d'Ulm, 75230 Paris CEDEX 05, France^(a)

⁽⁴⁾ Department of Physics, Technion-Israel Institute of Technology, 32000 Haifa, Israel

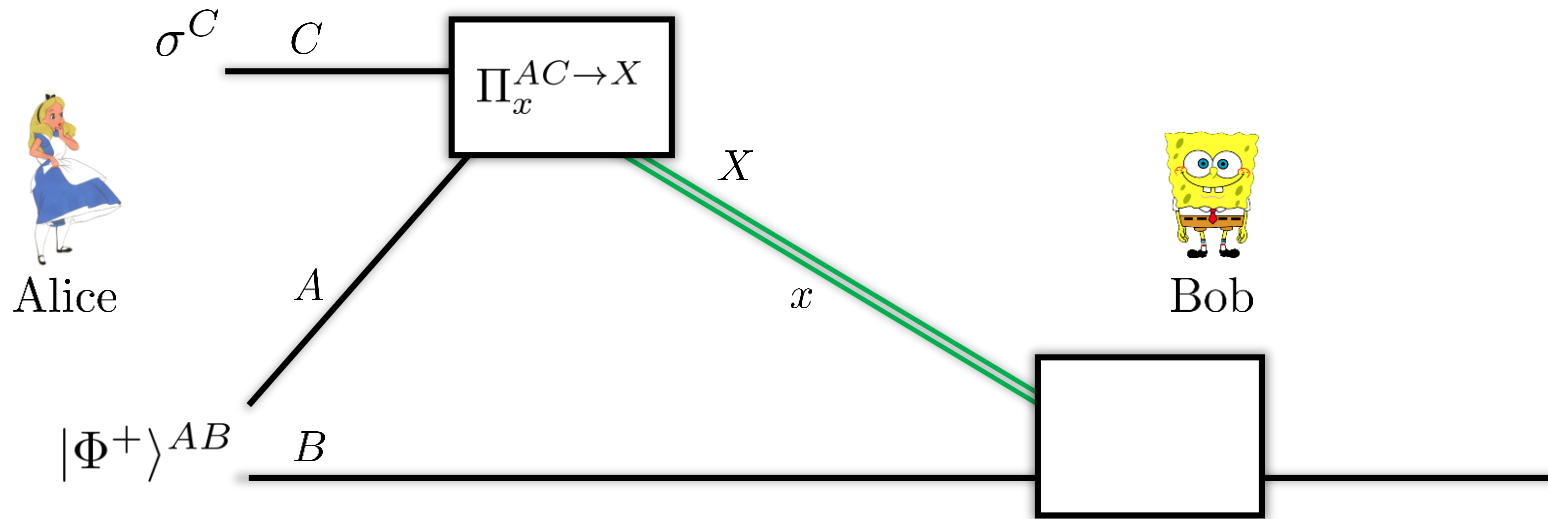
⁽⁵⁾ Department of Physics, Williams College, Williamstown, Massachusetts 01267

(Received 2 December 1992)

Quantum teleportation

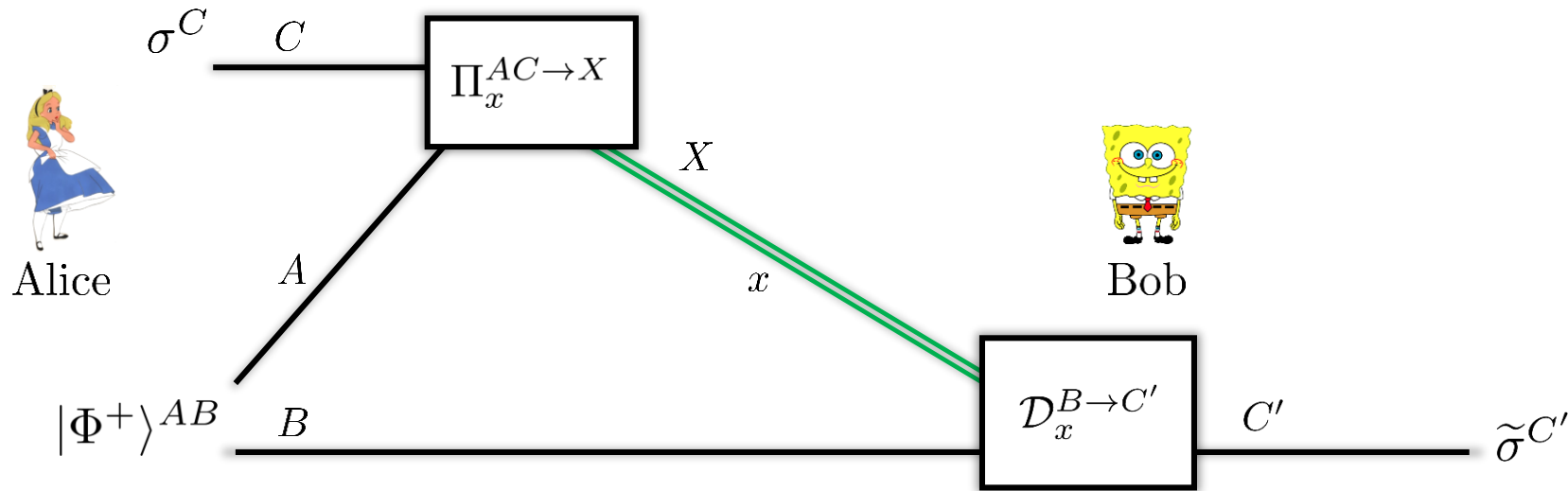


Quantum teleportation



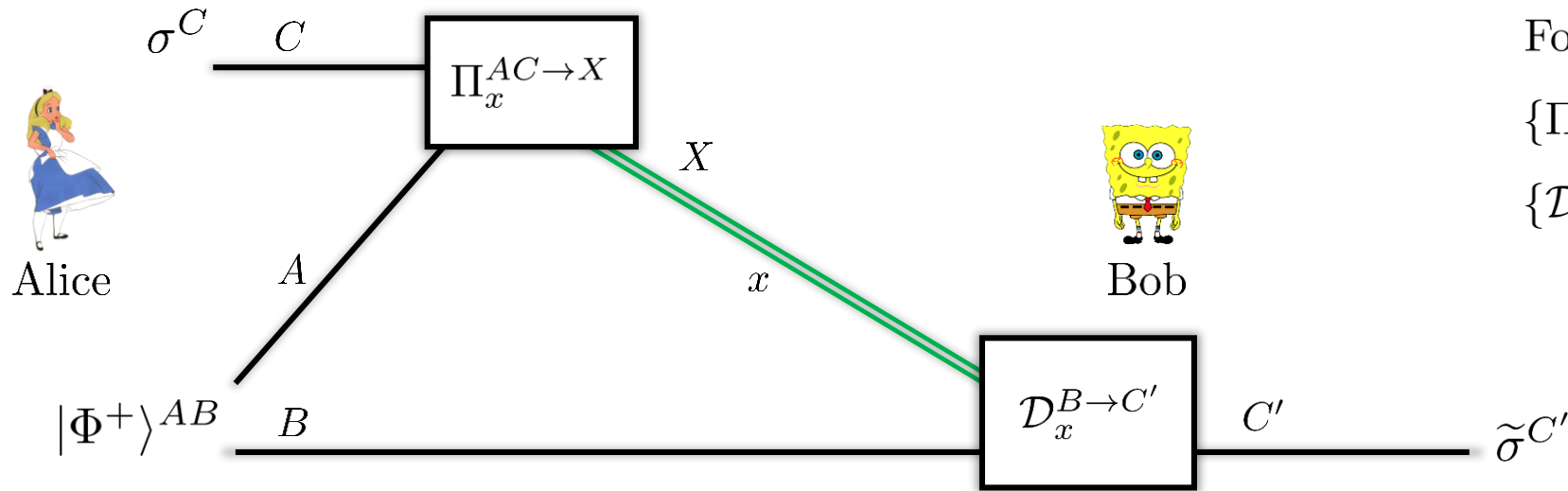
- Alice applies a qq-to-c encoder, which is a positive-operator valued-measure (POVM) $\{\Pi_x^{AC \rightarrow X}\}_x$

Quantum teleportation



- Alice applies a q-q-to-c encoder, which is a positive-operator valued-measure (POVM) $\{\Pi_x^{AC \rightarrow X}\}_x$
- Bob applies a family of q-to-q decoders, which is a family of completely-positive trace-preserving (CPTP) maps $\{\mathcal{D}_x^{B \rightarrow C'}\}_x$

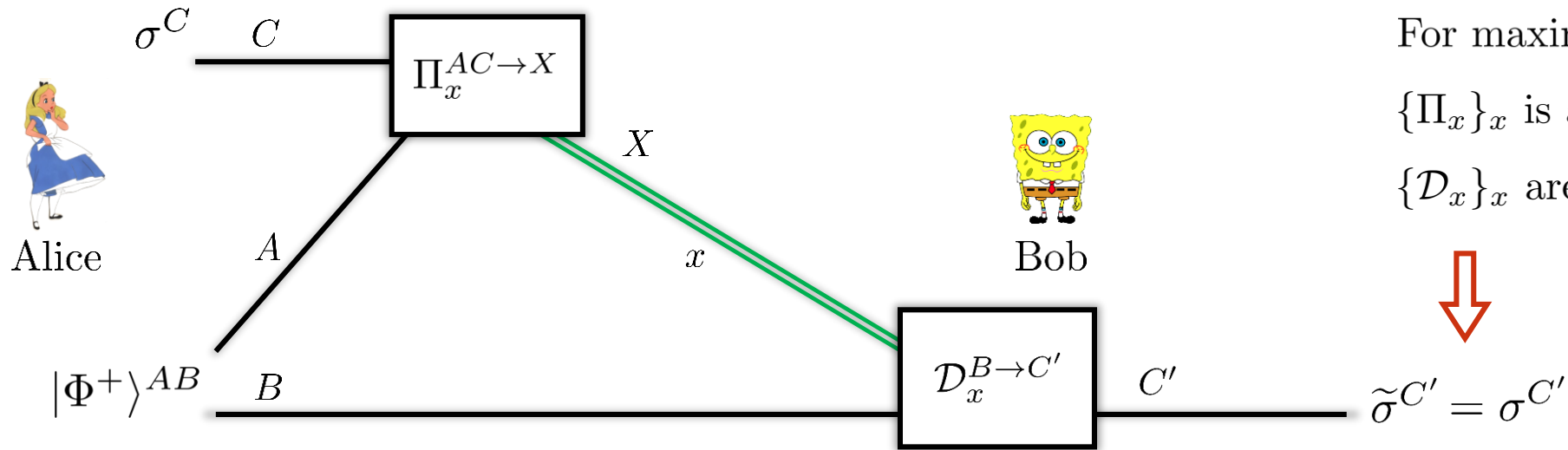
Quantum teleportation



For maximally entangled $|\Phi^+\rangle$,
 $\{\Pi_x\}_x$ is a Bell-state measurement
 $\{\mathcal{D}_x\}_x$ are Pauli channels

- Alice applies a q-q-to-c encoder, which is a positive-operator valued-measure (POVM) $\{\Pi_x^{AC \rightarrow X}\}_x$
- Bob applies a family of q-to-q decoders, which is a family of completely-positive trace-preserving (CPTP) maps $\{\mathcal{D}_x^{B \rightarrow C'}\}_x$

Quantum teleportation



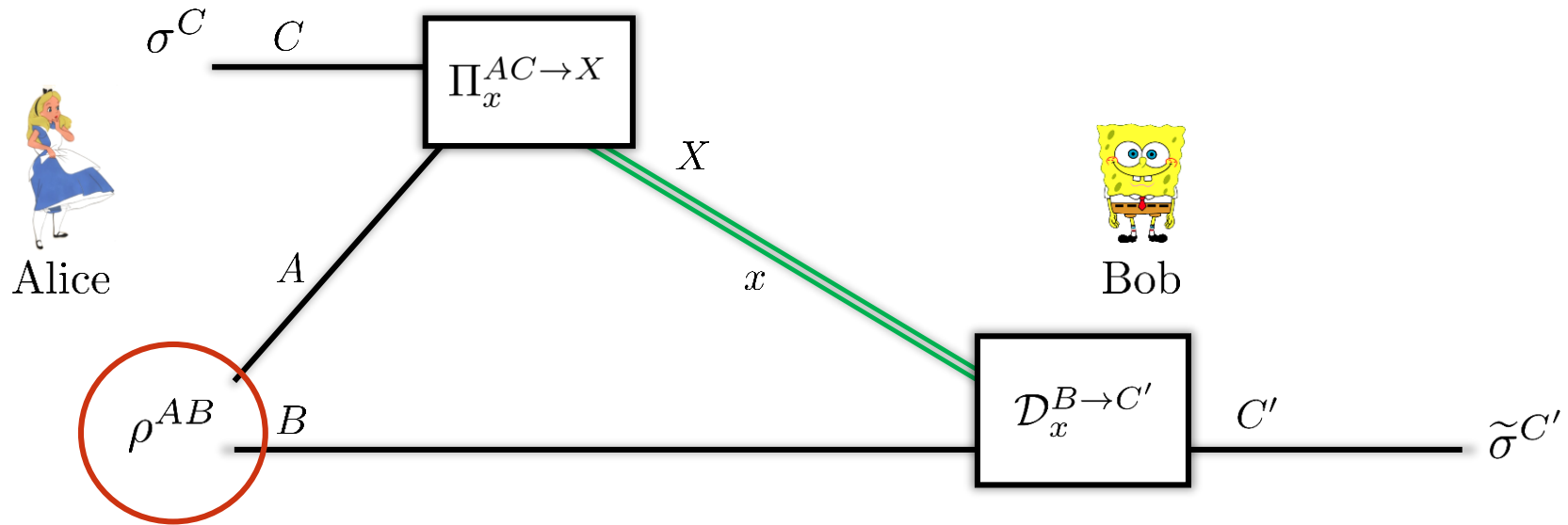
For maximally entangled $|\Phi^+\rangle$,
 $\{\Pi_x\}_x$ is a Bell-state measurement
 $\{\mathcal{D}_x\}_x$ are Pauli channels



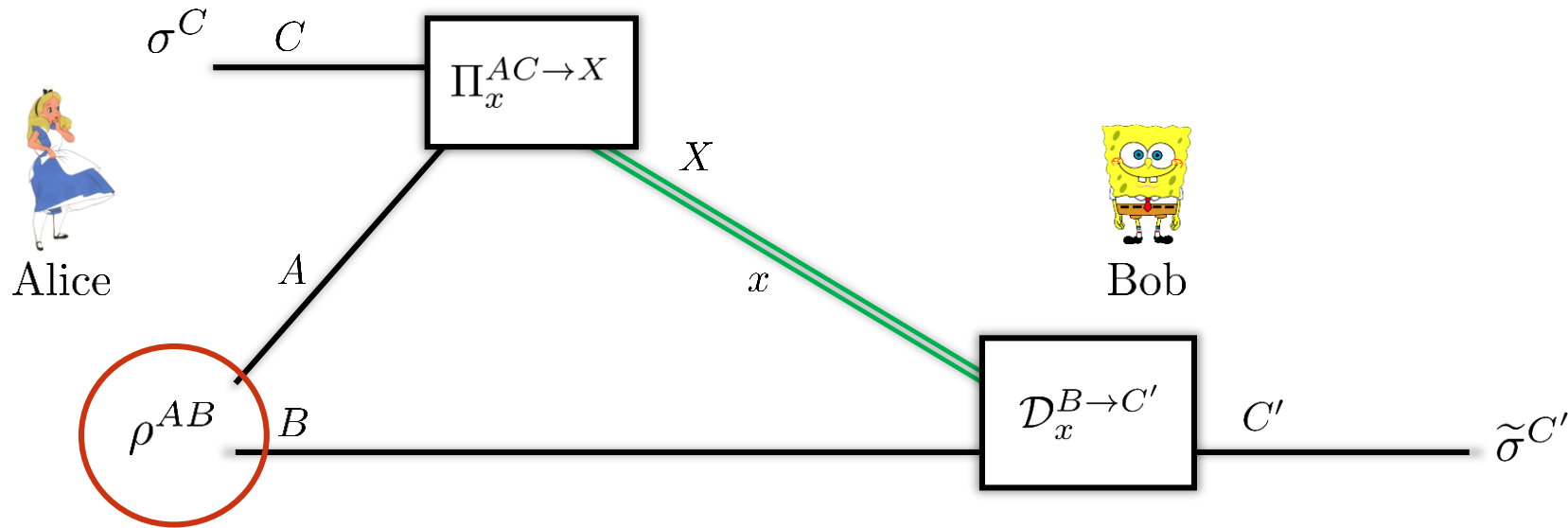
$$\tilde{\sigma}^{C'} = \sigma^{C'}$$

- Alice applies a q-q-to-c encoder, which is a positive-operator valued-measure (POVM) $\{\Pi_x^{AC \rightarrow X}\}_x$
- Bob applies a family of q-to-q decoders, which is a family of completely-positive trace-preserving (CPTP) maps $\{\mathcal{D}_x^{B \rightarrow C'}\}_x$

Quantum teleportation with noisy entanglement



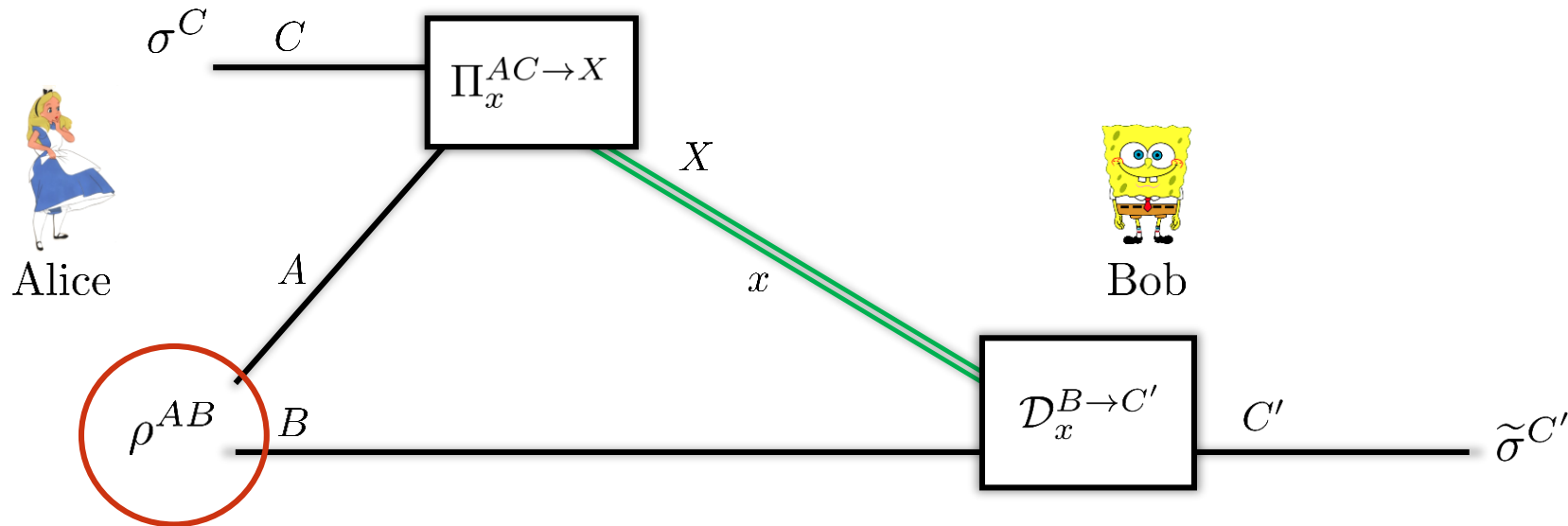
Quantum teleportation with noisy entanglement



- Every bipartite state ρ^{AB} , encoding POVM $\{\Pi_x^{AC}\}$, and decoding maps $\{\mathcal{D}_x^{B \rightarrow C'}\}_x$ define a $|C|$ -dimensional **one-way teleportation protocol** $(\rho, \{\Pi_x\}, \{\mathcal{D}_x\})$ for channel $\Lambda : C \rightarrow C'$,

$$\sigma \mapsto \Lambda(\sigma) = \sum_x \mathcal{D}_x^{B \rightarrow C'} \left(\text{Tr}_{AC} [(\mathbb{I}^B \otimes \Pi_x^{AC})(\rho^{AB} \otimes \sigma^C)] \right).$$

Quantum teleportation with noisy entanglement



- Every bipartite state ρ^{AB} , encoding POVM $\{\Pi_x^{AC}\}$, and decoding maps $\{\mathcal{D}_x^{B \rightarrow C'}\}_x$ define a $|C|$ -dimensional **one-way teleportation protocol** $(\rho, \{\Pi_x\}, \{\mathcal{D}_x\})$ for channel $\Lambda : C \rightarrow C'$,

$$\sigma \mapsto \Lambda(\sigma) = \sum_x \mathcal{D}_x^{B \rightarrow C'} \left(\text{Tr}_{AC} [(\mathbb{I}^B \otimes \Pi_x^{AC})(\rho^{AB} \otimes \sigma^C)] \right).$$

- How good of a quantum communication resource is this channel?

How good are general teleportation channels?

PHYSICAL REVIEW A

VOLUME 60, NUMBER 3

SEPTEMBER 1999

General teleportation channel, singlet fraction, and quasidistillation

Michał Horodecki*

Institute of Theoretical Physics and Astrophysics, University of Gdańsk, 80-952 Gdańsk, Poland

Paweł Horodecki†

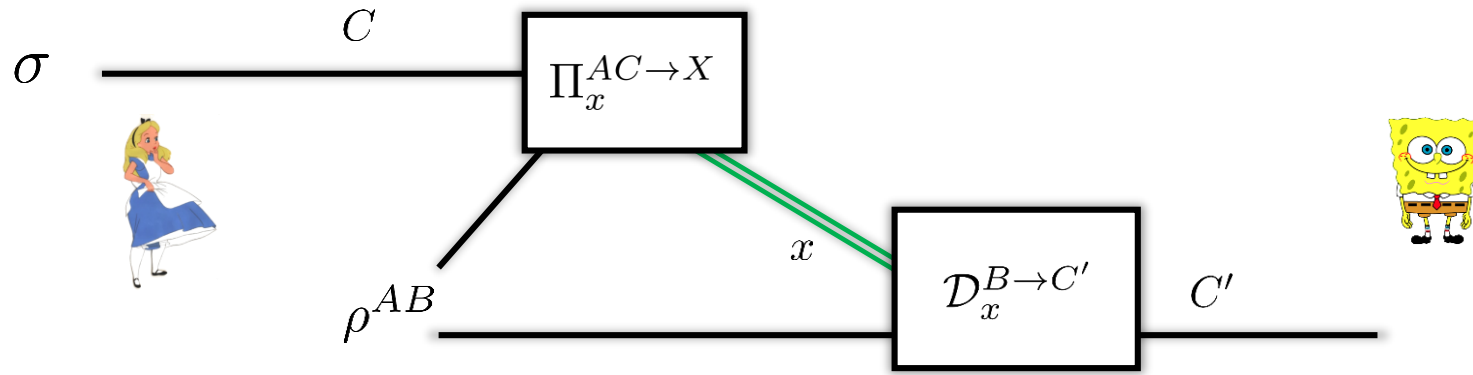
Faculty of Applied Physics and Mathematics, Technical University of Gdańsk, 80-952 Gdańsk, Poland

Ryszard Horodecki‡

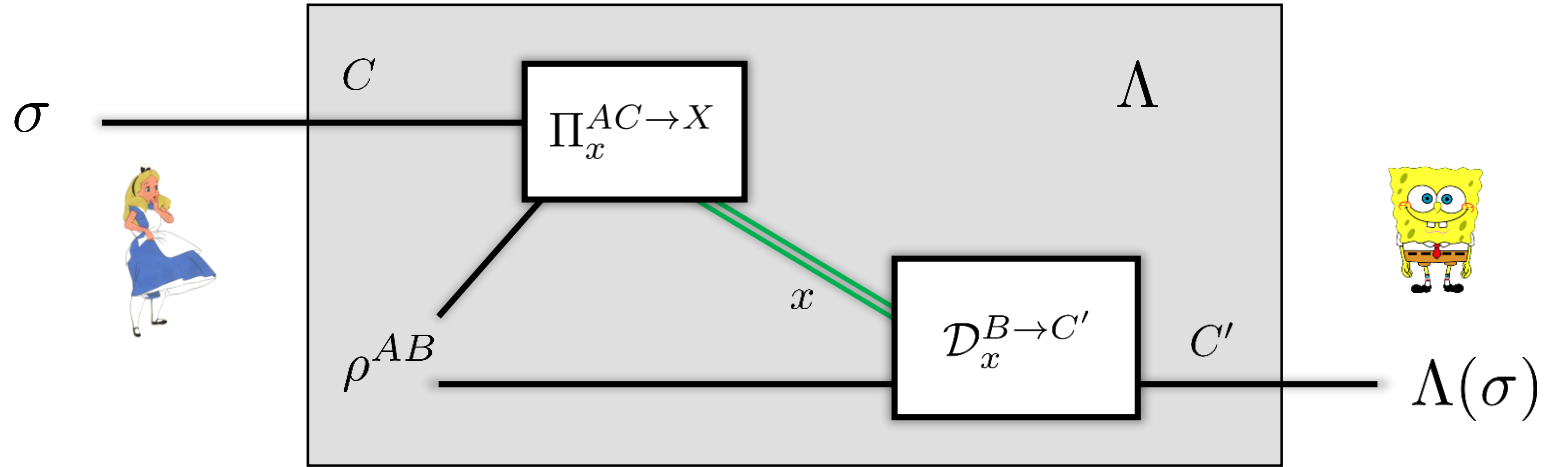
Institute of Theoretical Physics and Astrophysics, University of Gdańsk, 80-952 Gdańsk, Poland

(Received 21 August 1998)

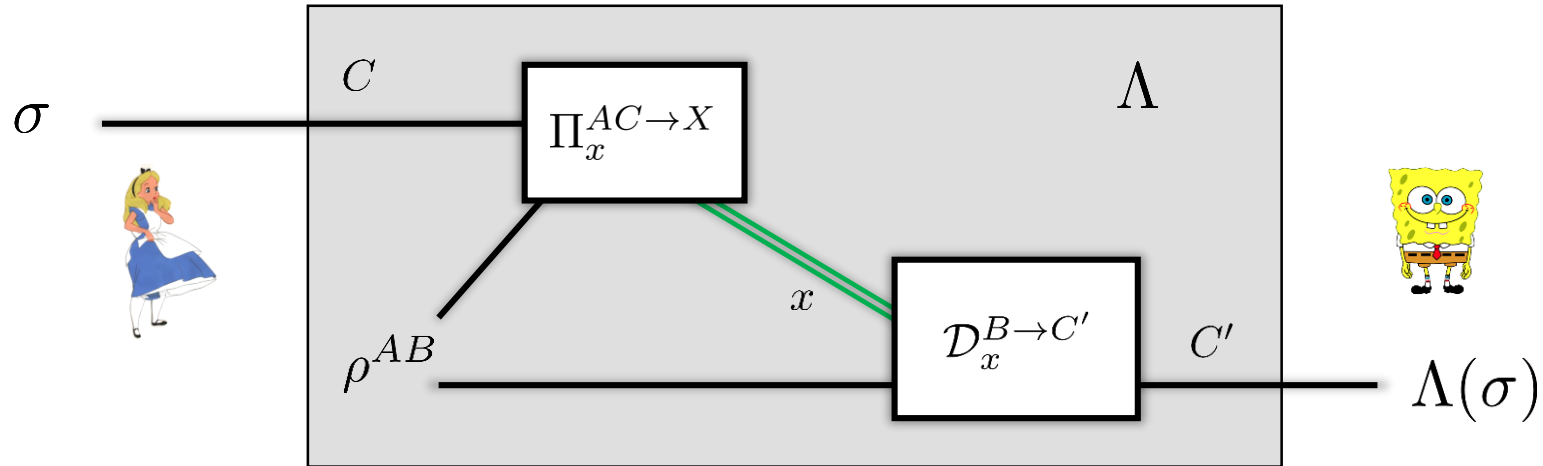
How good are general teleportation channels?



How good are general teleportation channels?



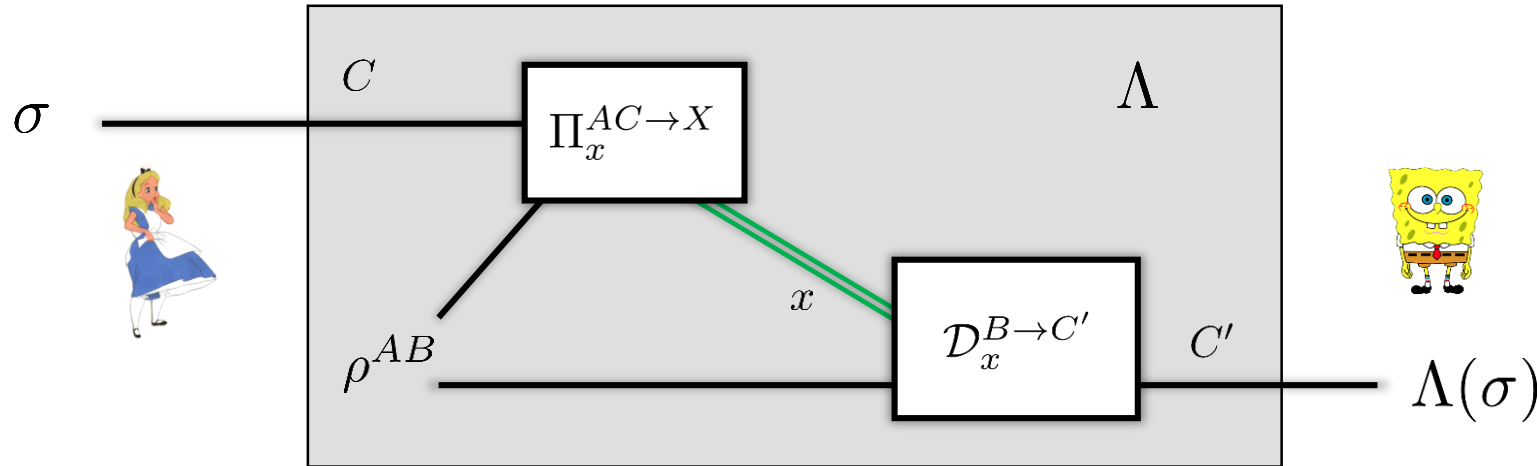
How good are general teleportation channels?



- The traditional measure used to assess the quality of a teleportation protocol is the **teleportation fidelity**:

$$f(\Lambda) = \int d\psi \langle \psi | \Lambda(\psi) | \psi \rangle.$$

How good are general teleportation channels?

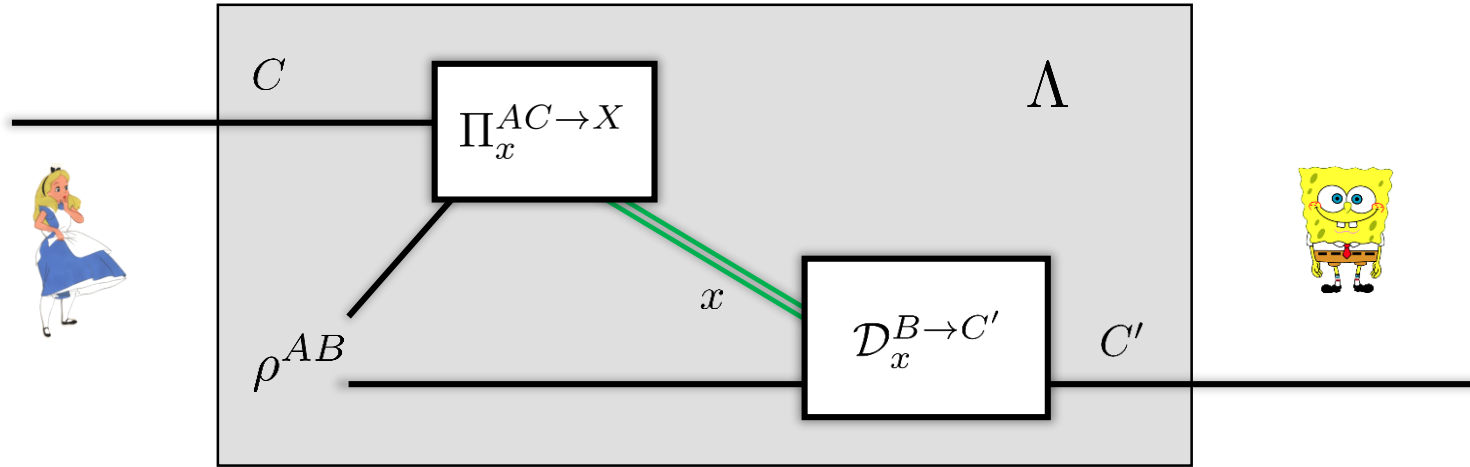


- The traditional measure used to assess the quality of a teleportation protocol is the **teleportation fidelity**:

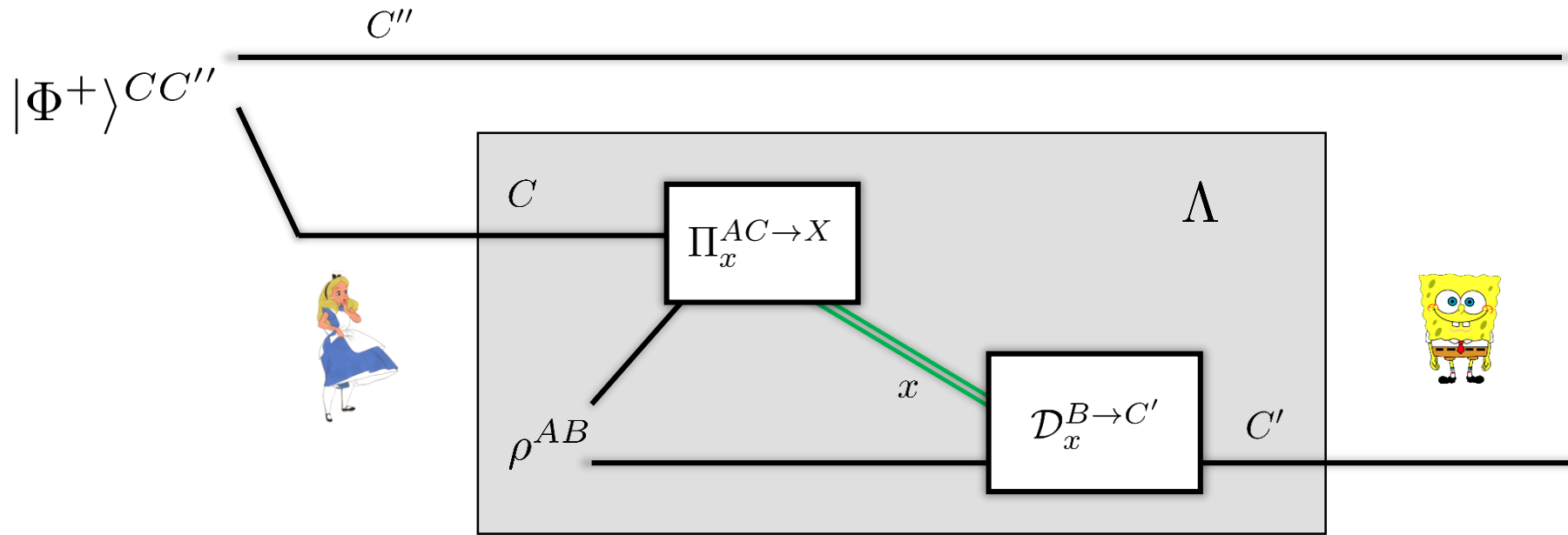
$$f(\Lambda) = \int d\psi \langle \psi | \Lambda(\psi) | \psi \rangle.$$

- This is the fidelity of transmission averaged over all input pure states $|\psi\rangle$.

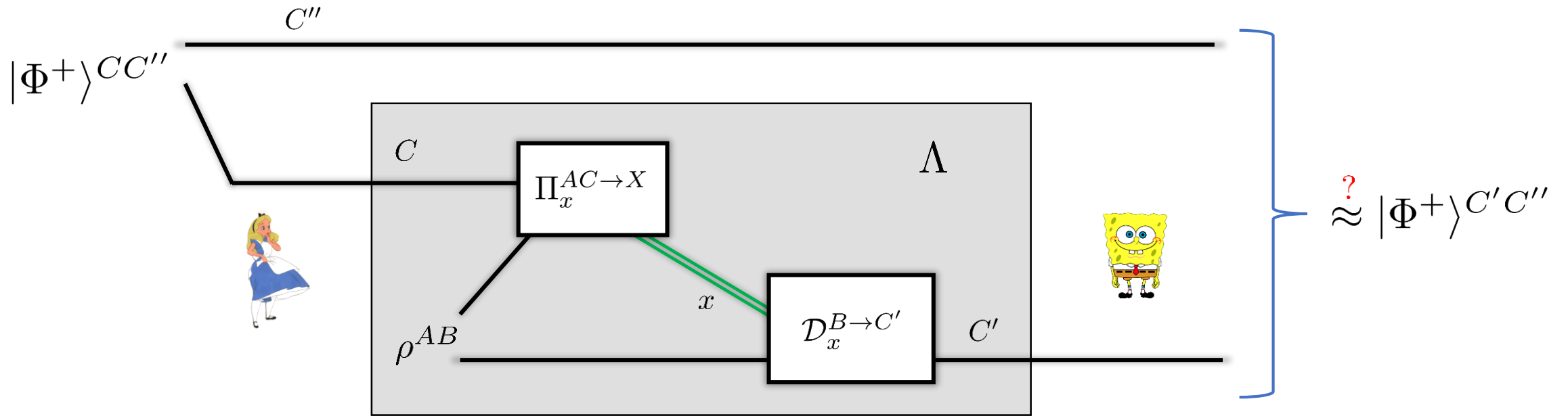
How good are general teleportation channels?



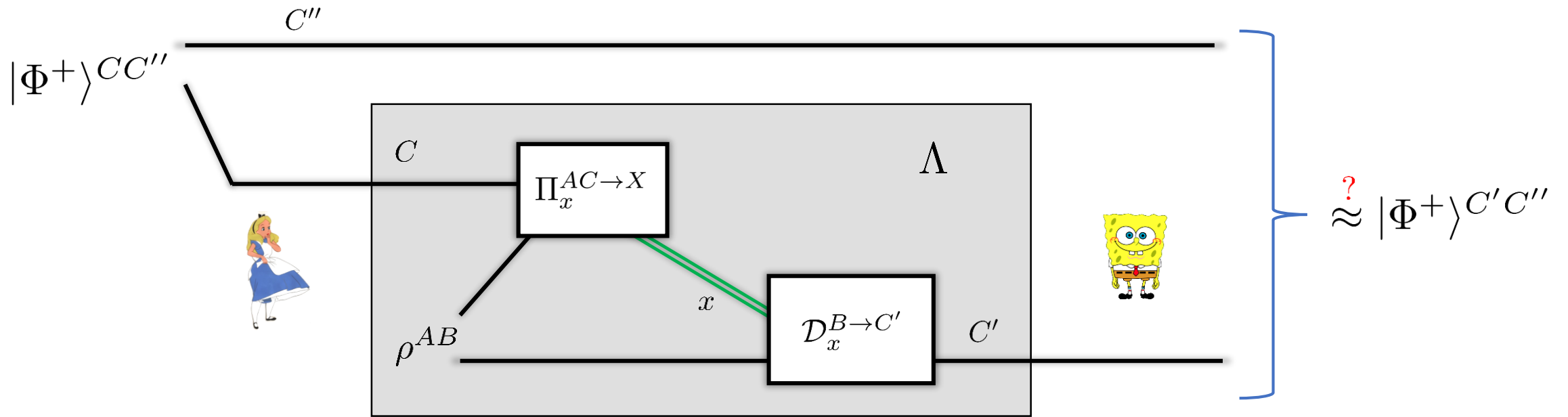
How good are general teleportation channels?



How good are general teleportation channels?



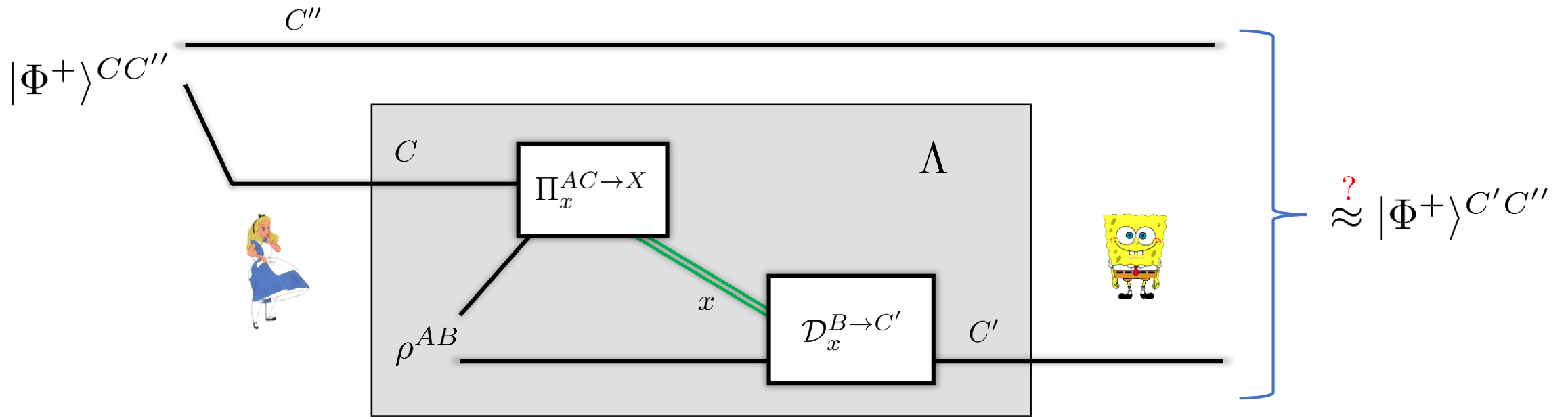
How good are general teleportation channels?



- An alternative measure is the **teleportation entanglement fidelity**.

$$F(\Lambda) = \langle \Phi^+ | \Lambda^{C \rightarrow C'} \otimes \text{id}^{C''} (\Phi^{+CC''}) | \Phi^+ \rangle_{C'C''}.$$

How good are general teleportation channels?



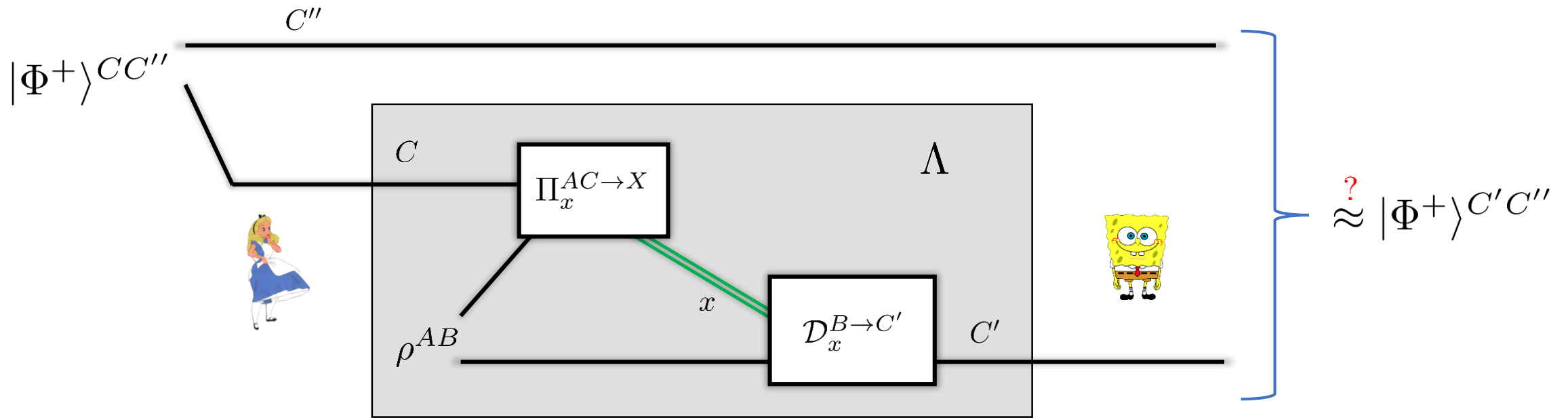
- An alternative measure is the **teleportation entanglement fidelity**.

$$F(\Lambda) = \langle \Phi^+ | \Lambda^{C \rightarrow C'} \otimes \text{id}^{C''} (\Phi^{+CC''}) | \Phi^+ \rangle_{C'C''}.$$

Theorem [Horodecki *et al.* PRA '99]

$$f(\Lambda) = \frac{F(\Lambda)|C| + 1}{|C| + 1}.$$

How good are general teleportation channels?



- An alternative measure is the **teleportation entanglement fidelity**.

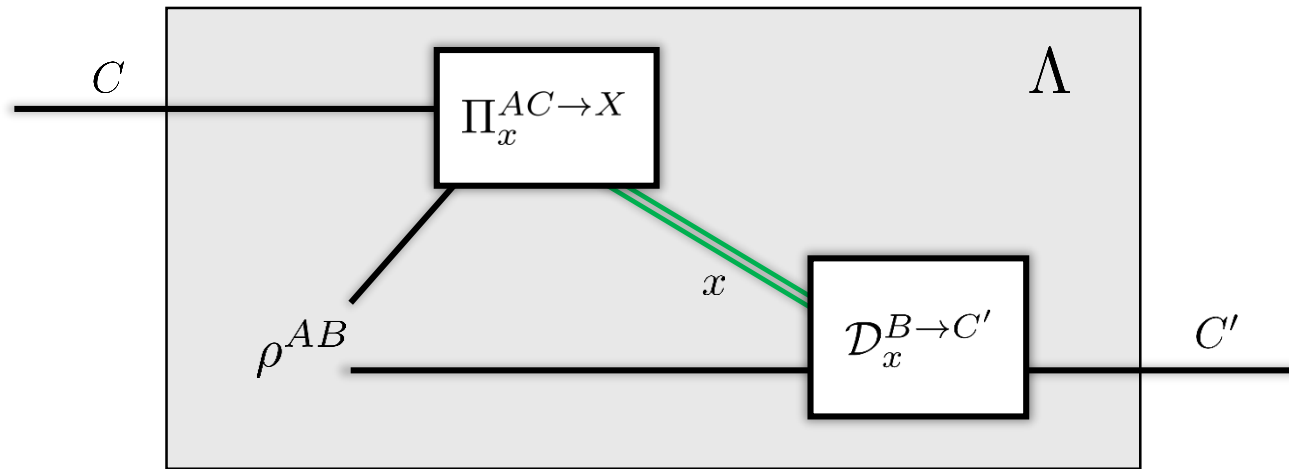
$$F(\Lambda) = \langle \Phi^+ | \Lambda^{C \rightarrow C'} \otimes \text{id}^{C''} (|\Phi^+\rangle_{CC''}) | \Phi^+ \rangle_{C'C''}.$$

Theorem [Horodecki *et al.* PRA '99]

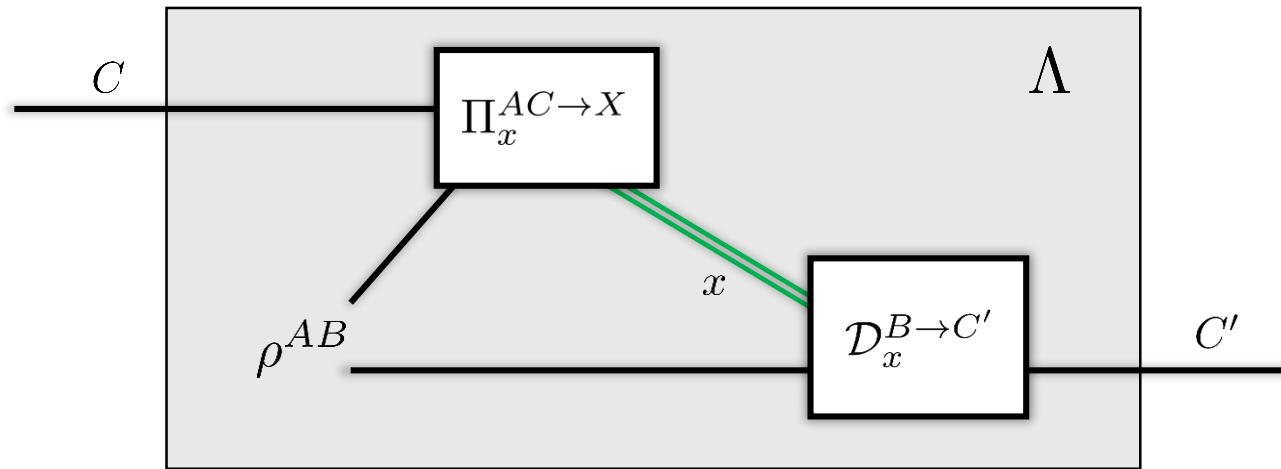
$$f(\Lambda) = \frac{F(\Lambda)|C| + 1}{|C| + 1}.$$

- **Our interest:** How else to interpret the teleportation entanglement fidelity $F(\Lambda)$?

General one-way teleportation protocols



General one-way teleportation protocols

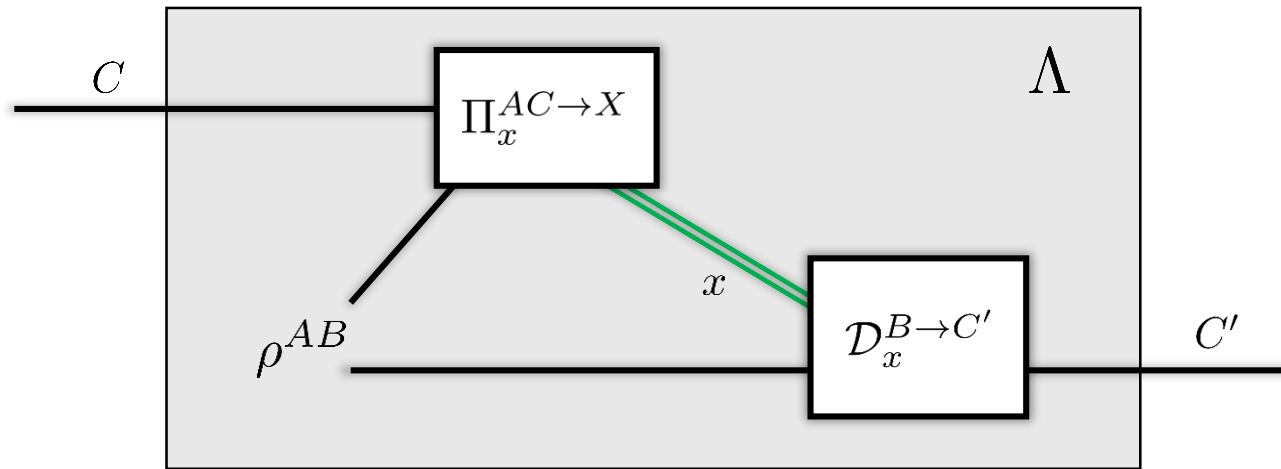


Chitambar, FL, arXiv:2302.14798

Lemma: Let $(\rho^{AB}, \{\Pi_x\}, \{\mathcal{D}_x\})$ define a one-way teleportation protocol. Then the teleportation entanglement fidelity is given by

$$F = \frac{N}{|C|^2} \left(\frac{1}{N} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}] \right), \quad \text{where } \omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB}).$$

General one-way teleportation protocols



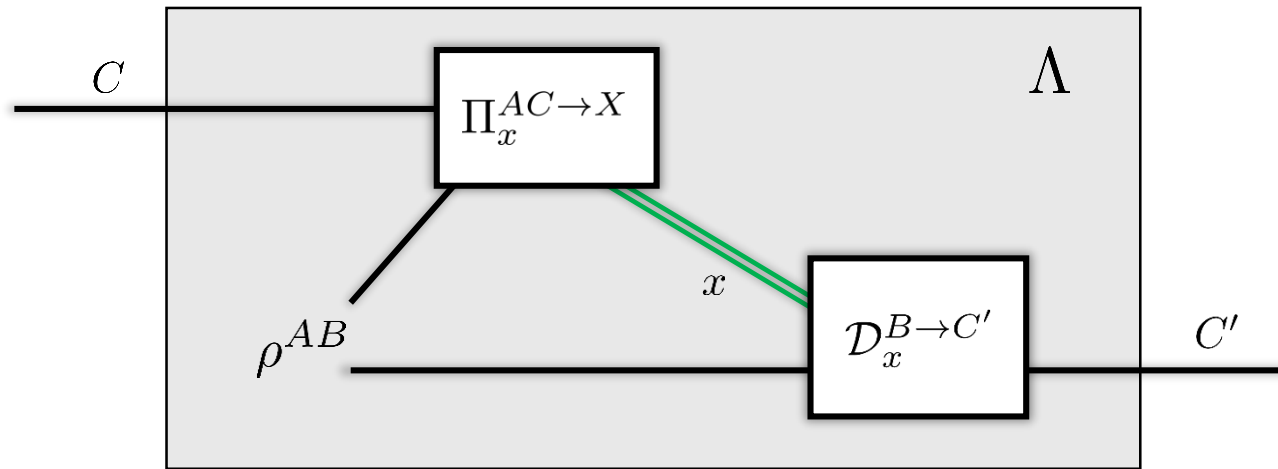
Chitambar, FL, arXiv:2302.14798

Lemma: Let $(\rho^{AB}, \{\Pi_x\}, \{\mathcal{D}_x\})$ define a one-way teleportation protocol. Then the teleportation entanglement fidelity is given by

$$F = \frac{N}{|C|^2} \left(\frac{1}{N} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}] \right), \quad \text{where } \omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB}).$$

Interpretation: The entanglement fidelity of every one-way teleportation protocol is equivalent to the success probability of a corresponding quantum state discrimination protocol.

General one-way teleportation protocols



Chitambar, FL, arXiv:2302.14798

Lemma: Let $(\rho^{AB}, \{\Pi_x\}, \{\mathcal{D}_x\})$ define a one-way teleportation protocol. Then the teleportation entanglement fidelity is given by

$$F = \frac{N}{|C|^2} \left(\frac{1}{N} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}] \right), \quad \text{where } \omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB}).$$

Interpretation: The entanglement fidelity of every one-way teleportation protocol is equivalent to the success probability of a corresponding quantum state discrimination protocol.

This **generalizes** a similar formula for port-based teleportation protocols to **arbitrary teleportation protocols**.

Ishizaka, Hiroshima, PRL 101, 240501 (2008)

The teleportation fidelity of a state

- How useful is a given quantum state ρ^{AB} for teleportation?

The teleportation fidelity of a state

- How useful is a given quantum state ρ^{AB} for teleportation?

Definition:

The $|C|$ -dimensional **teleportation fidelity** of state ρ^{AB} is defined as the maximum entanglement fidelity among all $|C|$ -dimensional teleportation protocols using ρ^{AB} :

$$F(\rho^{AB}; |C|) := \max_{\Lambda} F(\Lambda)$$

The teleportation fidelity of a state

- How useful is a given quantum state ρ^{AB} for teleportation?

Definition:

The $|C|$ -dimensional **teleportation fidelity** of state ρ^{AB} is defined as the maximum entanglement fidelity among all $|C|$ -dimensional teleportation protocols using ρ^{AB} :

$$F(\rho^{AB}; |C|) := \max_{\Lambda} F(\Lambda)$$

$$= \max_{\{\Pi_x\}_x, \{\mathcal{D}_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}], \quad \text{where } \omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB}).$$

The teleportation fidelity of a state

- How useful is a given quantum state ρ^{AB} for teleportation?

Definition:

The $|C|$ -dimensional **teleportation fidelity** of state ρ^{AB} is defined as the maximum entanglement fidelity among all $|C|$ -dimensional teleportation protocols using ρ^{AB} :

$$F(\rho^{AB}; |C|) := \max_{\Lambda} F(\Lambda)$$

$$= \max_{\{\Pi_x\}_x, \{\mathcal{D}_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}], \quad \text{where } \omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB}).$$

- For the classical preparation channel $\mathcal{D}_x(\cdot) = \text{Tr}(\cdot)|x\rangle\langle x|$ and measurement $\Pi_x = \mathbb{I} \otimes |x\rangle\langle x|$, we have

$$\text{Tr}[\Pi_x \omega_x] = \text{Tr}[(\mathbb{I} \otimes |x\rangle\langle x|)(\rho^A \otimes |x\rangle\langle x|^C)] = 1.$$

The teleportation fidelity of a state

- How useful is a given quantum state ρ^{AB} for teleportation?

Definition:

The $|C|$ -dimensional **teleportation fidelity** of state ρ^{AB} is defined as the maximum entanglement fidelity among all $|C|$ -dimensional teleportation protocols using ρ^{AB} :

$$F(\rho^{AB}; |C|) := \max_{\Lambda} F(\Lambda)$$

$$= \max_{\{\Pi_x\}_x, \{\mathcal{D}_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}], \quad \text{where } \omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB}).$$

- For the classical preparation channel $\mathcal{D}_x(\cdot) = \text{Tr}(\cdot)|x\rangle\langle x|$ and measurement $\Pi_x = \mathbb{I} \otimes |x\rangle\langle x|$, we have

$$\text{Tr}[\Pi_x \omega_x] = \text{Tr}[(\mathbb{I} \otimes |x\rangle\langle x|)(\rho^A \otimes |x\rangle\langle x|^C)] = 1.$$



$$F(\rho^{AB}; |C|) \geq \frac{1}{|C|}.$$

The teleportation fidelity of a state

- How useful is a given quantum state ρ^{AB} for teleportation?

Definition:

The $|C|$ -dimensional **teleportation fidelity** of state ρ^{AB} is defined as the maximum entanglement fidelity among all $|C|$ -dimensional teleportation protocols using ρ^{AB} :

$$F(\rho^{AB}; |C|) := \max_{\Lambda} F(\Lambda)$$

$$= \max_{\{\Pi_x\}_x, \{\mathcal{D}_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}], \quad \text{where } \omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB}).$$

- For the classical preparation channel $\mathcal{D}_x(\cdot) = \text{Tr}(\cdot)|x\rangle\langle x|$ and measurement $\Pi_x = \mathbb{I} \otimes |x\rangle\langle x|$, we have

$$\text{Tr}[\Pi_x \omega_x] = \text{Tr}[(\mathbb{I} \otimes |x\rangle\langle x|)(\rho^A \otimes |x\rangle\langle x|^C)] = 1.$$



$$F(\rho^{AB}; |C|) \geq \frac{1}{|C|}.$$

- This lower bound is called the **classical teleportation threshold**.

The teleportation fidelity of a state

- How useful is a given quantum state ρ^{AB} for teleportation?

Definition:

The $|C|$ -dimensional **teleportation fidelity** of state ρ^{AB} is defined as the maximum entanglement fidelity among all $|C|$ -dimensional teleportation protocols using ρ^{AB} :

$$F(\rho^{AB}; |C|) := \max_{\Lambda} F(\Lambda)$$

$$= \max_{\{\Pi_x\}_x, \{\mathcal{D}_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}], \quad \text{where } \omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB}).$$

- For the classical preparation channel $\mathcal{D}_x(\cdot) = \text{Tr}(\cdot)|x\rangle\langle x|$ and measurement $\Pi_x = \mathbb{I} \otimes |x\rangle\langle x|$, we have

$$\text{Tr}[\Pi_x \omega_x] = \text{Tr}[(\mathbb{I} \otimes |x\rangle\langle x|)(\rho^A \otimes |x\rangle\langle x|^C)] = 1.$$



$$F(\rho^{AB}; |C|) \geq \frac{1}{|C|}.$$

- This lower bound is called the **classical teleportation threshold**.

Fundamental question: When is this lower bound tight?

The reduction criterion and the classical threshold

The reduction criterion and the classical threshold

- Recall that a bipartite state ρ^{AB} is said to satisfy the **reduction criterion** if

$$\rho^A \otimes \mathbb{I}^B - \rho^{AB} \geq 0.$$

Horodecki & Horodecki PRA **59** 4206 (1999)

The reduction criterion and the classical threshold

- Recall that a bipartite state ρ^{AB} is said to satisfy the **reduction criterion** if

$$\rho^A \otimes \mathbb{I}^B - \rho^{AB} \geq 0.$$

Horodecki & Horodecki PRA **59** 4206 (1999)

- (i) Every separable (non-entangled) state satisfies the reduction criterion.

The reduction criterion and the classical threshold

- Recall that a bipartite state ρ^{AB} is said to satisfy the **reduction criterion** if

$$\rho^A \otimes \mathbb{I}^B - \rho^{AB} \geq 0.$$

Horodecki & Horodecki PRA **59** 4206 (1999)

- (i) Every separable (non-entangled) state satisfies the reduction criterion.
- (ii) Every state violating the reduction criterion is distillable.

The reduction criterion and the classical threshold

- Recall that a bipartite state ρ^{AB} is said to satisfy the **reduction criterion** if

$$\rho^A \otimes \mathbb{I}^B - \rho^{AB} \geq 0.$$

Horodecki & Horodecki PRA **59** 4206 (1999)

- (i) Every separable (non-entangled) state satisfies the reduction criterion.
- (ii) Every state violating the reduction criterion is distillable.

 If ρ^{AB} is not distillable, then $\omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion.

The reduction criterion and the classical threshold

- Recall that a bipartite state ρ^{AB} is said to satisfy the **reduction criterion** if

$$\rho^A \otimes \mathbb{I}^B - \rho^{AB} \geq 0.$$

Horodecki & Horodecki PRA **59** 4206 (1999)

- (i) Every separable (non-entangled) state satisfies the reduction criterion.
- (ii) Every state violating the reduction criterion is distillable.

⇒ If ρ^{AB} is not distillable, then $\omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion.

$$\Rightarrow \omega_x^{AC} \leq \rho^A \otimes \mathbb{I}^C \quad \text{for all } x$$

The reduction criterion and the classical threshold

- Recall that a bipartite state ρ^{AB} is said to satisfy the **reduction criterion** if

$$\rho^A \otimes \mathbb{I}^B - \rho^{AB} \geq 0.$$

Horodecki & Horodecki PRA **59** 4206 (1999)

- (i) Every separable (non-entangled) state satisfies the reduction criterion.
- (ii) Every state violating the reduction criterion is distillable.

⇒ If ρ^{AB} is not distillable, then $\omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion.

$$\Rightarrow \omega_x^{AC} \leq \rho^A \otimes \mathbb{I}^C \quad \text{for all } x$$

$$\Rightarrow \max_{\{\Pi_x\}_x, \{\mathcal{D}_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}] \leq \max_{\{\Pi_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} (\rho^A \otimes \mathbb{I}^C)] = \frac{1}{|C|}.$$

The reduction criterion and the classical threshold

- Recall that a bipartite state ρ^{AB} is said to satisfy the **reduction criterion** if

$$\rho^A \otimes \mathbb{I}^B - \rho^{AB} \geq 0.$$

Horodecki & Horodecki PRA **59** 4206 (1999)

- (i) Every separable (non-entangled) state satisfies the reduction criterion.
- (ii) Every state violating the reduction criterion is distillable.

⇒ If ρ^{AB} is not distillable, then $\omega_x = \text{id}^A \otimes \mathcal{D}_x^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion.

$$\Rightarrow \omega_x^{AC} \leq \rho^A \otimes \mathbb{I}^C \quad \text{for all } x$$

$$\Rightarrow \max_{\{\Pi_x\}_x, \{\mathcal{D}_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} \omega_x^{AC}] \leq \max_{\{\Pi_x\}_x} \frac{1}{|C|^2} \sum_{x=1}^N \text{Tr}[\Pi_x^{AC} (\rho^A \otimes \mathbb{I}^C)] = \frac{1}{|C|}.$$

Corollary: Bound entangled states cannot exceed the (one-way) teleportation classical threshold.

Horodecki *et al.* PRA **60** 1888 (1999)

The reduction criterion and the classical threshold

Theorem: A bipartite state ρ^{AB} is useful for $|C|$ -dimensional teleportation (i.e., exceed the classical fidelity threshold $|C|^{-1}$) iff there exists a channel $\mathcal{E}^{B \rightarrow C}$ such that $\omega^{AC} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ violates the reduction criterion.

Chitambar, FL, arXiv:2302.14798

The reduction criterion and the classical threshold

Theorem: A bipartite state ρ^{AB} is useful for $|C|$ -dimensional teleportation (i.e., exceed the classical fidelity threshold $|C|^{-1}$) iff there exists a channel $\mathcal{E}^{B \rightarrow C}$ such that $\omega^{AC} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ violates the reduction criterion.

Chitambar, FL, arXiv:2302.14798

- Compare with the classic result of the Horodeckis:

Theorem: A bipartite state ρ^{AB} is useful for $|C|$ -dimensional teleportation (i.e., exceed the classical fidelity threshold $|C|^{-1}$) iff there exists a one-way LOCC map \mathcal{L} such that $\mathcal{L}(\rho)$ has a singlet fraction exceeding $|C|^{-1}$:

$$\langle \Phi^+ | \mathcal{L}(\rho) | \Phi^+ \rangle^{C\tilde{C}} > \frac{1}{|C|}.$$

Horodecki *et al.* PRA **60** 1888 (1999)

The reduction criterion and the classical threshold

Theorem: A bipartite state ρ^{AB} is useful for $|C|$ -dimensional teleportation (i.e., exceed the classical fidelity threshold $|C|^{-1}$) iff there exists a channel $\mathcal{E}^{B \rightarrow C}$ such that $\omega^{AC} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ violates the reduction criterion.

Chitambar, FL, arXiv:2302.14798

- Compare with the classic result of the Horodeckis:

Theorem: A bipartite state ρ^{AB} is useful for $|C|$ -dimensional teleportation (i.e., exceed the classical fidelity threshold $|C|^{-1}$) iff there exists a one-way LOCC map \mathcal{L} such that $\mathcal{L}(\rho)$ has a singlet fraction exceeding $|C|^{-1}$:

$$\langle \Phi^+ | \mathcal{L}(\rho) | \Phi^+ \rangle^{C\tilde{C}} > \frac{1}{|C|}.$$

Horodecki *et al.* PRA **60** 1888 (1999)

- Our work simplifies the condition for non-classical teleportation fidelity from an optimization over all one-way LOCC maps to an optimization over just local maps.

Example: qutrit Werner states

Example: qutrit Werner states

- Consider the $3 \otimes 3$ family of Werner states:

$$\rho_\lambda^{AB} = \frac{1}{24}[(3-\lambda)\mathbb{I}_3^A \otimes \mathbb{I}_3^B + (3\lambda-1)\mathbb{F}_3^{AB}], \quad \text{where } \mathbb{F} \text{ is the SWAP operator.}$$

Example: qutrit Werner states

- Consider the $3 \otimes 3$ family of Werner states:

$$\rho_\lambda^{AB} = \frac{1}{24}[(3-\lambda)\mathbb{I}_3^A \otimes \mathbb{I}_3^B + (3\lambda-1)\mathbb{F}_3^{AB}], \quad \text{where } \mathbb{F} \text{ is the SWAP operator.}$$

- Every ρ_λ satisfies the reduction criterion, but ρ_λ is entangled iff $\lambda < 0$.

Example: qutrit Werner states

- Consider the $3 \otimes 3$ family of Werner states:

$$\rho_\lambda^{AB} = \frac{1}{24}[(3-\lambda)\mathbb{I}_3^A \otimes \mathbb{I}_3^B + (3\lambda-1)\mathbb{F}_3^{AB}], \quad \text{where } \mathbb{F} \text{ is the SWAP operator.}$$

- Every ρ_λ satisfies the reduction criterion, but ρ_λ is entangled iff $\lambda < 0$.
- However, consider the CPTP on Bob's system that collapses it to a two-dimensional subspace,

$$\mathcal{E}(\rho) = (|0\rangle\langle 0| + |1\rangle\langle 1|)\rho(|0\rangle\langle 0| + |1\rangle\langle 1|) + |0\rangle\langle 2|\rho|2\rangle\langle 0|.$$

Example: qutrit Werner states

- Consider the $3 \otimes 3$ family of Werner states:

$$\rho_\lambda^{AB} = \frac{1}{24}[(3-\lambda)\mathbb{I}_3^A \otimes \mathbb{I}_3^B + (3\lambda-1)\mathbb{F}_3^{AB}], \quad \text{where } \mathbb{F} \text{ is the SWAP operator.}$$

- Every ρ_λ satisfies the reduction criterion, but ρ_λ is entangled iff $\lambda < 0$.
- However, consider the CPTP on Bob's system that collapses it to a two-dimensional subspace,

$$\mathcal{E}(\rho) = (|0\rangle\langle 0| + |1\rangle\langle 1|)\rho(|0\rangle\langle 0| + |1\rangle\langle 1|) + |0\rangle\langle 2|\rho|2\rangle\langle 0|.$$

- Easy calculation: $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho_\lambda^{AB})$ violates the reduction criterion for $-1 \leq \lambda < -3/7$.

Example: qutrit Werner states

- Consider the $3 \otimes 3$ family of Werner states:

$$\rho_\lambda^{AB} = \frac{1}{24}[(3-\lambda)\mathbb{I}_3^A \otimes \mathbb{I}_3^B + (3\lambda-1)\mathbb{F}_3^{AB}], \quad \text{where } \mathbb{F} \text{ is the SWAP operator.}$$

- Every ρ_λ satisfies the reduction criterion, but ρ_λ is entangled iff $\lambda < 0$.
- However, consider the CPTP on Bob's system that collapses it to a two-dimensional subspace,

$$\mathcal{E}(\rho) = (|0\rangle\langle 0| + |1\rangle\langle 1|)\rho(|0\rangle\langle 0| + |1\rangle\langle 1|) + |0\rangle\langle 2|\rho|2\rangle\langle 0|.$$

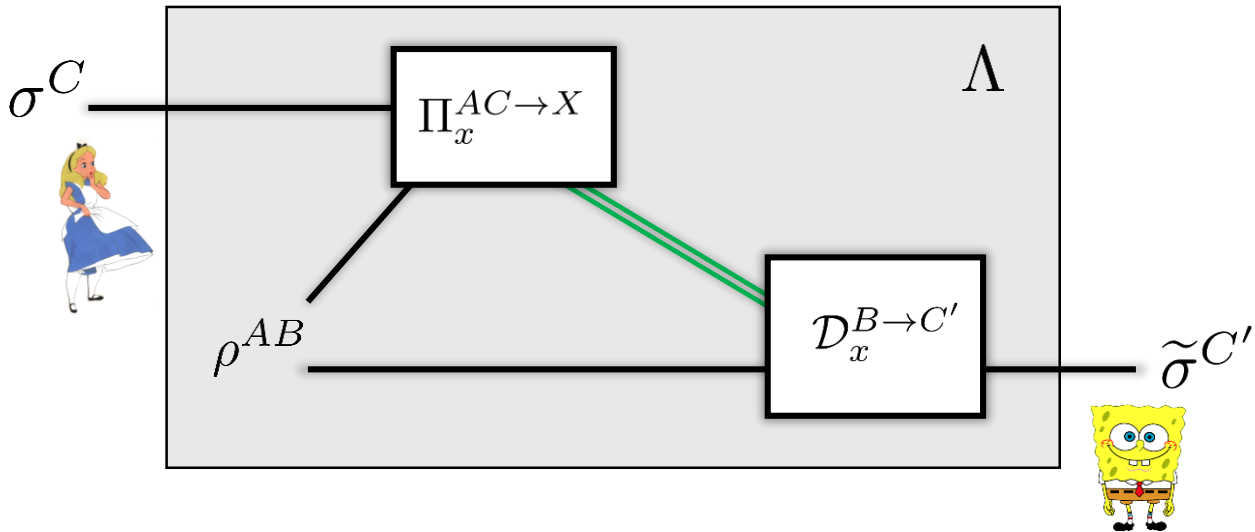
- Easy calculation: $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho_\lambda^{AB})$ violates the reduction criterion for $-1 \leq \lambda < -3/7$.

 Even states satisfying the reduction criterion can be useful for teleportation.

From teleportation to dense coding

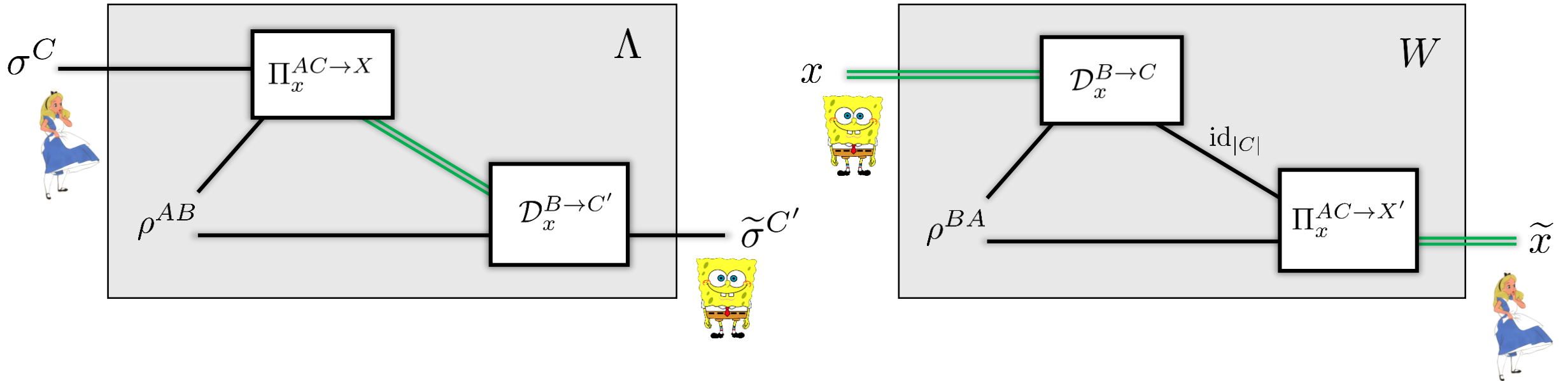
From teleportation to dense coding

ρ^{AB}	$\{\Pi_x\}_{x=1}^N$	$\{\mathcal{D}_x\}_{x=1}^N$
Alice-to-Bob teleportation	Quantum encoder	Quantum decoder



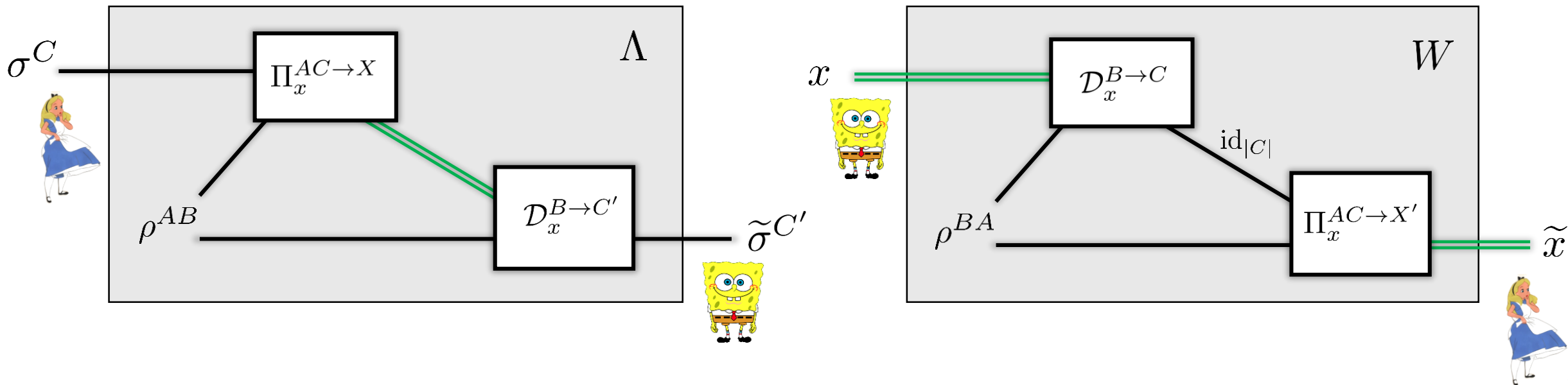
From teleportation to dense coding

ρ^{AB}	$\{\Pi_x\}_{x=1}^N$	$\{\mathcal{D}_x\}_{x=1}^N$
Alice-to-Bob teleportation	Quantum encoder	Quantum decoder



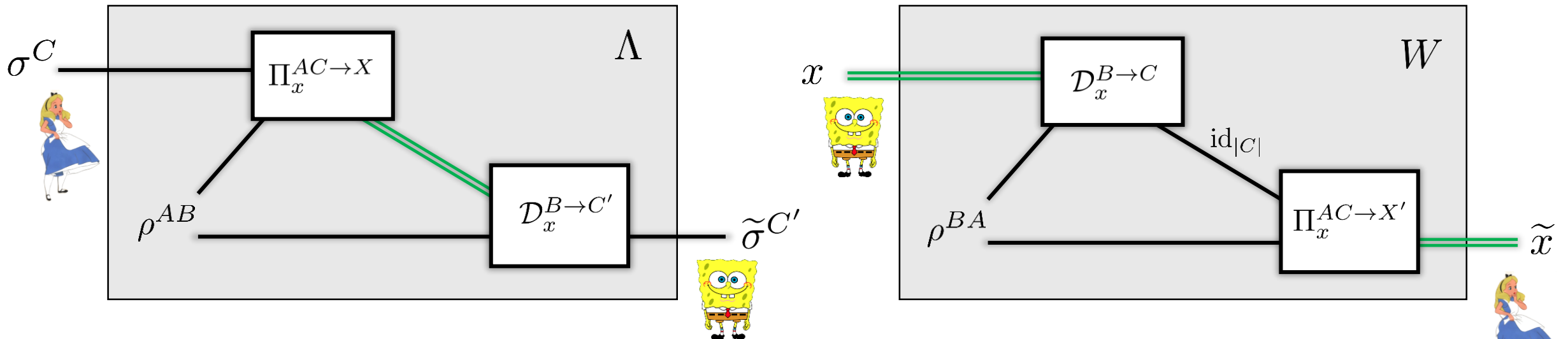
From teleportation to dense coding

ρ^{AB}	$\{\Pi_x\}_{x=1}^N$	$\{\mathcal{D}_x\}_{x=1}^N$
Alice-to-Bob teleportation	Quantum encoder	Quantum decoder
Bob-to-Alice dense coding	Classical decoder	Classical encoder



From teleportation to dense coding

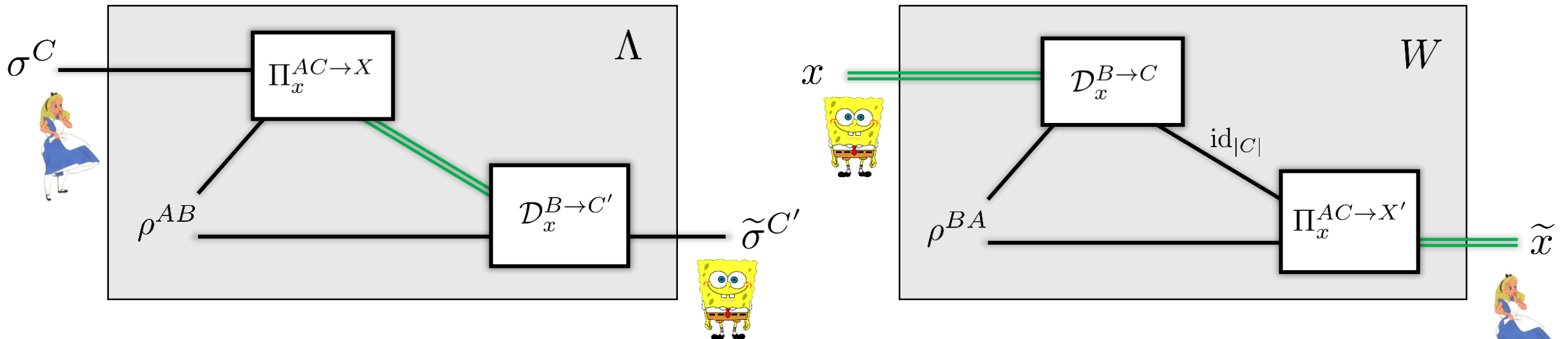
ρ^{AB}	$\{\Pi_x\}_{x=1}^N$	$\{\mathcal{D}_x\}_{x=1}^N$
Alice-to-Bob teleportation	Quantum encoder	Quantum decoder
Bob-to-Alice dense coding	Classical decoder	Classical encoder



- Every teleportation channel $\Lambda : C \rightarrow C'$ using ρ has a corresponding dense coding channel $W : \{1, \dots, N\} \rightarrow \{1, \dots, N\}$, and vice-versa.

From teleportation to dense coding

ρ^{AB}	$\{\Pi_x\}_{x=1}^N$	$\{\mathcal{D}_x\}_{x=1}^N$
Alice-to-Bob teleportation	Quantum encoder	Quantum decoder
Bob-to-Alice dense coding	Classical decoder	Classical encoder



- Every teleportation channel $\Lambda : C \rightarrow C'$ using ρ has a corresponding dense coding channel $W : \{1, \dots, N\} \rightarrow \{1, \dots, N\}$, and vice-versa.

- “Dense coding” corresponds to the situation when $N > |C|$.

How good are general dense coding channels?

How good are general dense coding channels?

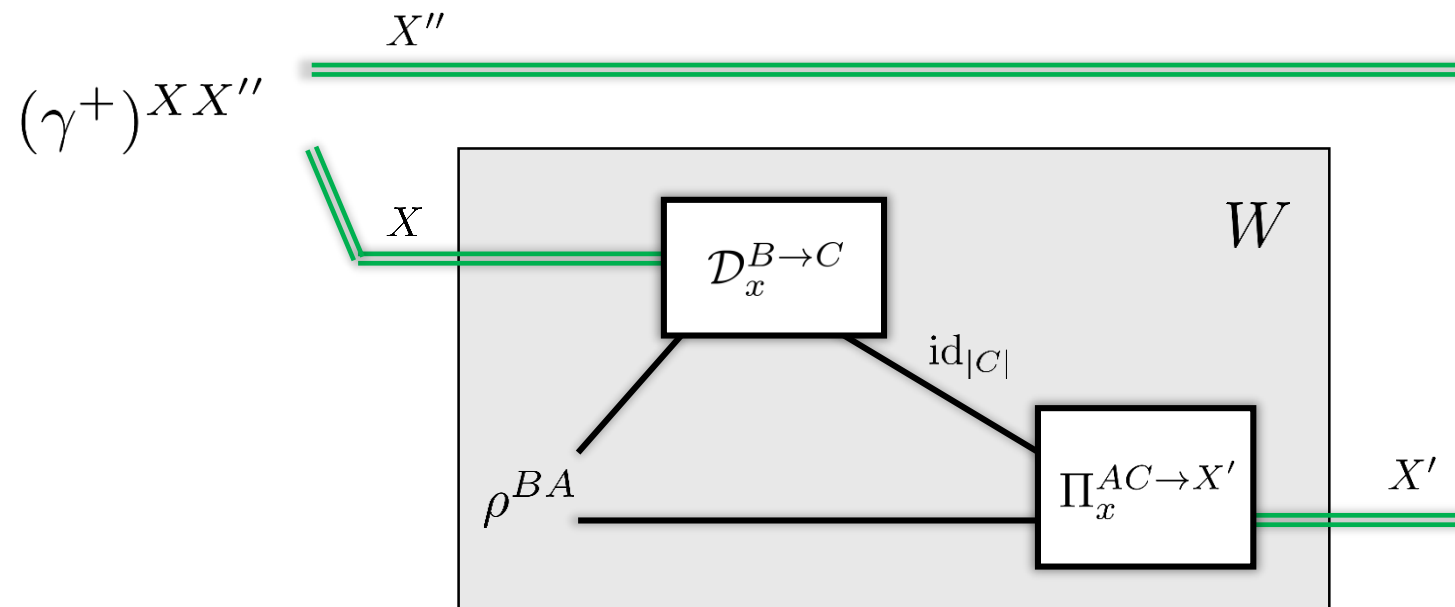
- Let $(\gamma^+)^{XX'} = \frac{1}{N} \sum_{x=1}^N |x\rangle\langle x|^X \otimes |x\rangle\langle x|^{X'}$ define an N -dimensional **maximally correlated** state.

How good are general dense coding channels?

- Let $(\gamma^+)^{XX'} = \frac{1}{N} \sum_{x=1}^N |x\rangle\langle x|^X \otimes |x\rangle\langle x|^{X'}$ define an N -dimensional **maximally correlated** state.
- This is the classical analog of the maximally entangled state $|\Phi^+\rangle^{CC'} = \frac{1}{\sqrt{|C|}} \sum_{i=1}^{|C|} |i\rangle_C \otimes |i\rangle_{C'}$.

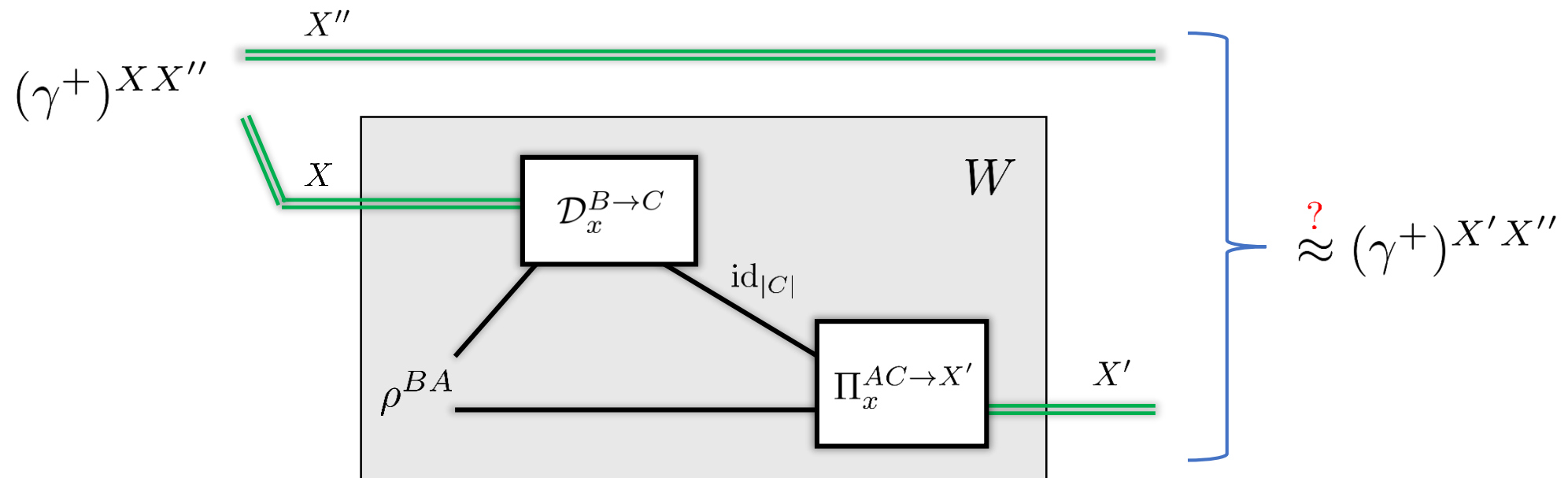
How good are general dense coding channels?

- Let $(\gamma^+)^{XX'} = \frac{1}{N} \sum_{x=1}^N |x\rangle\langle x|^X \otimes |x\rangle\langle x|^{X'}$ define an N -dimensional **maximally correlated** state.
- This is the classical analog of the maximally entangled state $|\Phi^+\rangle^{CC'} = \frac{1}{\sqrt{|C|}} \sum_{i=1}^{|C|} |i\rangle_C \otimes |i\rangle_{C'}$.



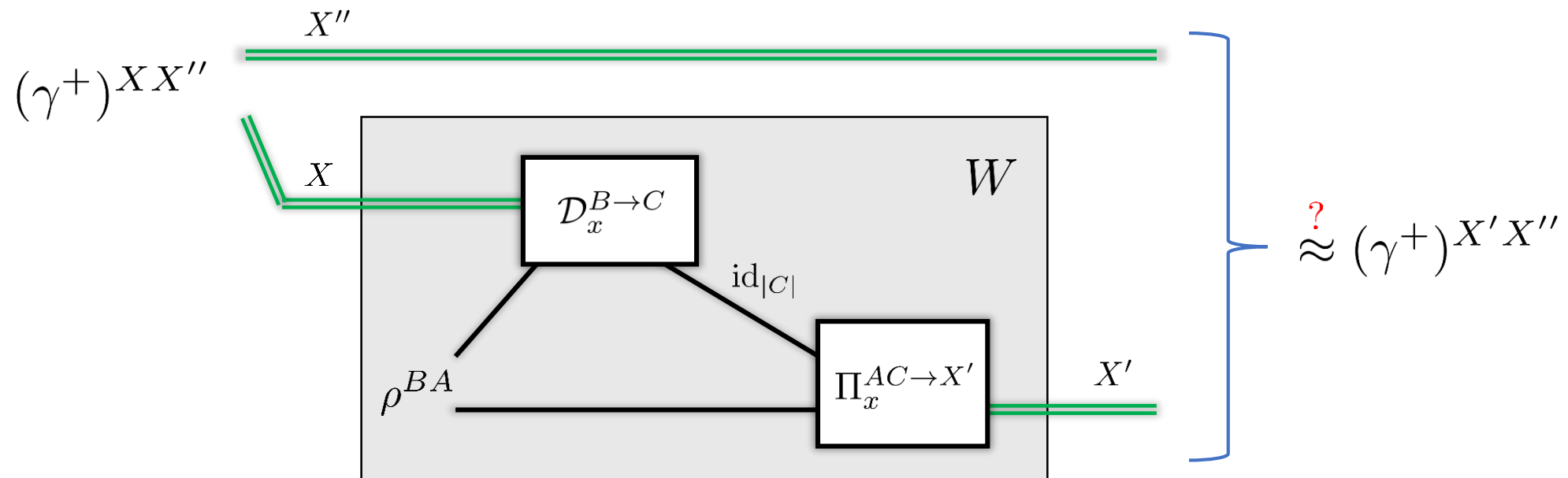
How good are general dense coding channels?

- Let $(\gamma^+)^{XX'} = \frac{1}{N} \sum_{x=1}^N |x\rangle\langle x|^X \otimes |x\rangle\langle x|^{X'}$ define an N -dimensional **maximally correlated** state.
- This is the classical analog of the maximally entangled state $|\Phi^+\rangle^{CC'} = \frac{1}{\sqrt{|C|}} \sum_{i=1}^{|C|} |i\rangle_C \otimes |i\rangle_{C'}$.



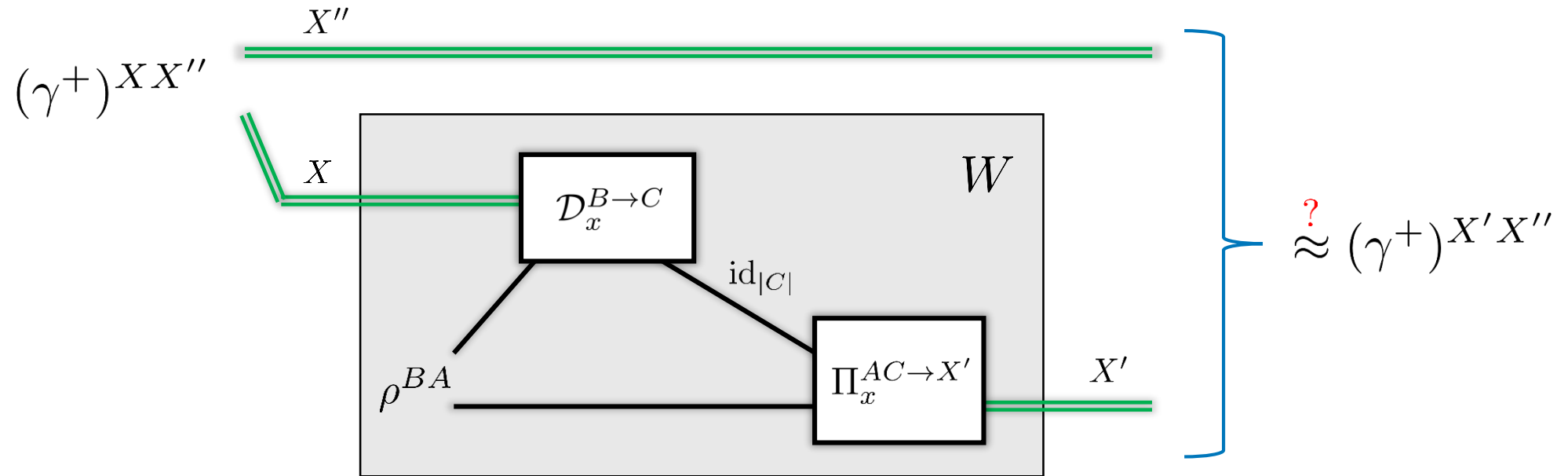
How good are general dense coding channels?

- Let $(\gamma^+)^{XX'} = \frac{1}{N} \sum_{x=1}^N |x\rangle\langle x|^X \otimes |x\rangle\langle x|^{X'}$ define an N -dimensional **maximally correlated** state.
- This is the classical analog of the maximally entangled state $|\Phi^+\rangle^{CC'} = \frac{1}{\sqrt{|C|}} \sum_{i=1}^{|C|} |i\rangle_C \otimes |i\rangle_{C'}$.

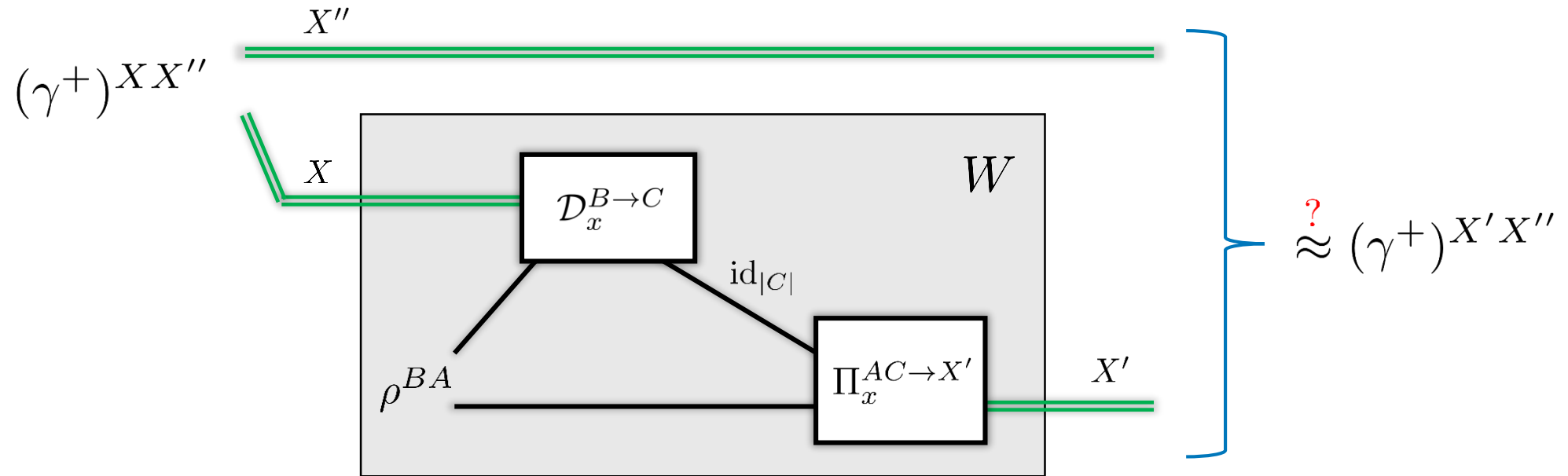


- One way to measure the quality of the classical channel is how well it preserves the maximally correlated (classical) state.

How good are general dense coding channels?



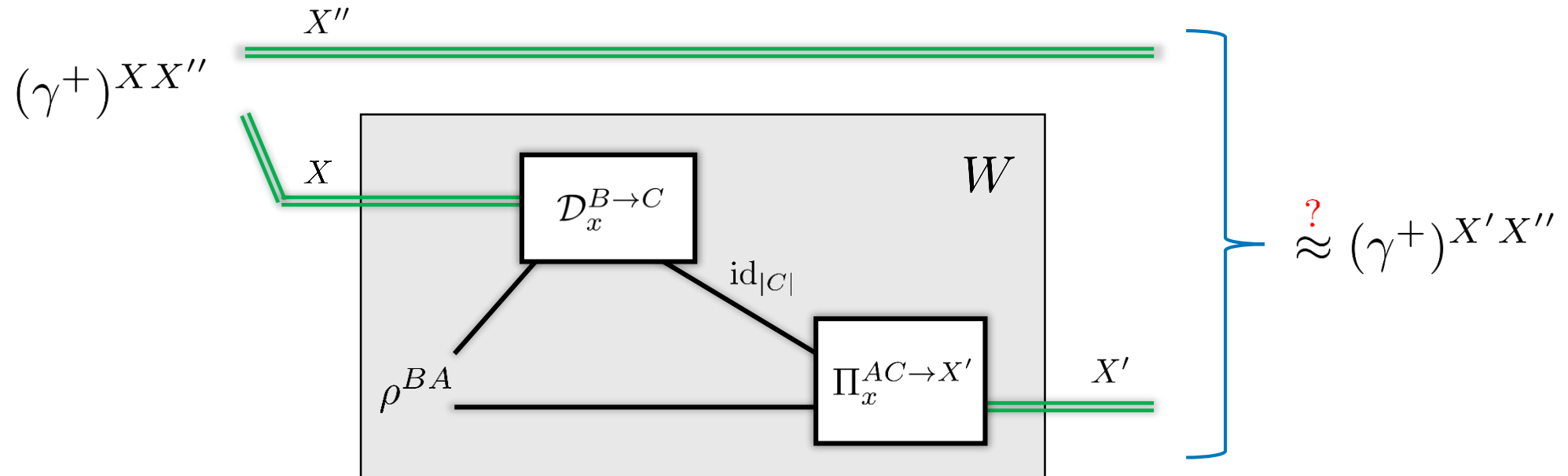
How good are general dense coding channels?



- Analogous to the entanglement fidelity of a quantum channel, we define the **classical correlation fidelity** of a channel $W: X \rightarrow X'$ as

$$F_{\text{cl}}(W) = \text{Tr}[(\gamma^+)^{X' X''} W^{X \rightarrow X'} \otimes \text{id}^{X''} (\gamma^+)^{X X''}]$$

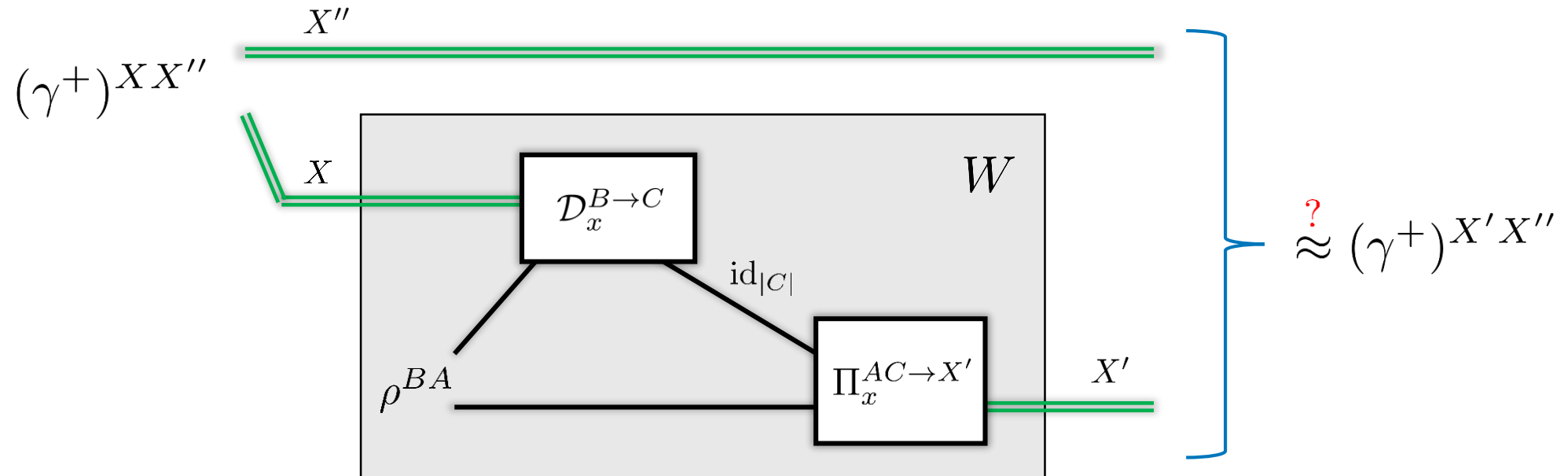
How good are general dense coding channels?



- Analogous to the entanglement fidelity of a quantum channel, we define the **classical correlation fidelity** of a channel $W: X \rightarrow X'$ as

$$\begin{aligned}
 F_{\text{cl}}(W) &= \text{Tr}[(\gamma^+)^{X'X''} W^{X \rightarrow X'} \otimes \text{id}^{X''} (\gamma^+)^{XX''}] \\
 &= \frac{1}{N} \sum_{x=1}^N W(x|x)
 \end{aligned}$$

How good are general dense coding channels?



- Analogous to the entanglement fidelity of a quantum channel, we define the **classical correlation fidelity** of a channel $W: X \rightarrow X'$ as

$$\begin{aligned}
 F_{\text{cl}}(W) &= \text{Tr}[(\gamma^+)^{X' X''} W^{X \rightarrow X'} \otimes \text{id}^{X''} (\gamma^+)^{X X''}] \\
 &= \frac{1}{N} \sum_{x=1}^N W(x|x)
 \end{aligned}$$

Interpretation: The classical correlation fidelity of channel W is equivalent to how well the channel transmits a randomly chosen message $x \in \{1, \dots, N\}$.

Teleportation and dense coding duality

- The classical threshold of the classical fidelity is $\min\{\frac{|C|}{N}, 1\}$.

Teleportation and dense coding duality

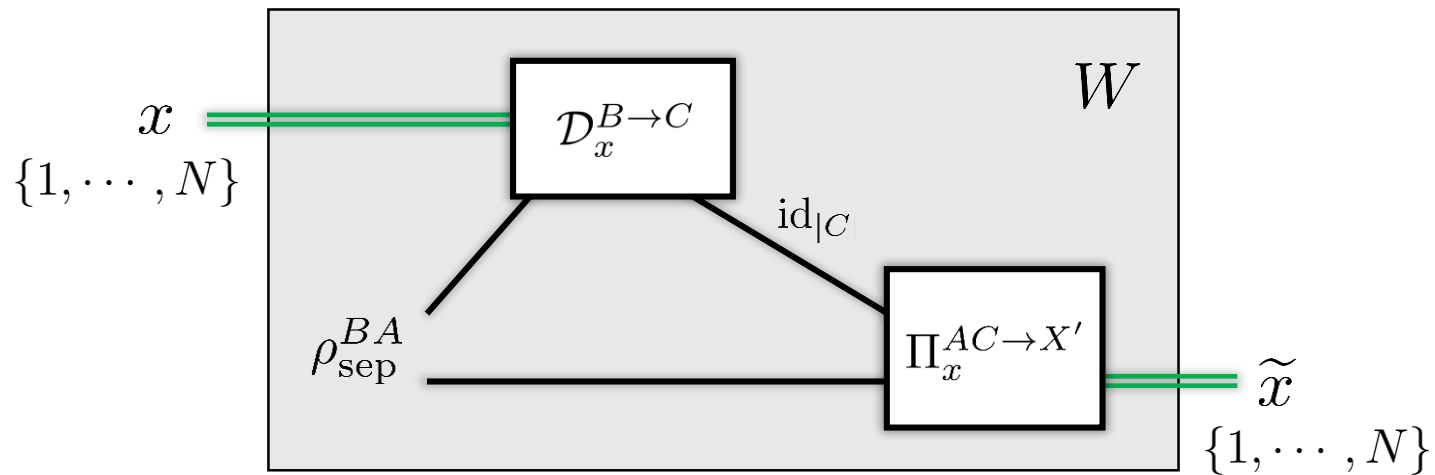
- The classical threshold of the classical fidelity is $\min\{\frac{|C|}{N}, 1\}$.

Interpretation: Using any non-entangled state ρ^{AB} and a $|C|$ -dimensional (noiseless) quantum channel, Alice can send Bob N classical messages with $p_{\text{succ}} \leq \min\{\frac{|C|}{N}, 1\}$.

Teleportation and dense coding duality

- The classical threshold of the classical fidelity is $\min\{\frac{|C|}{N}, 1\}$.

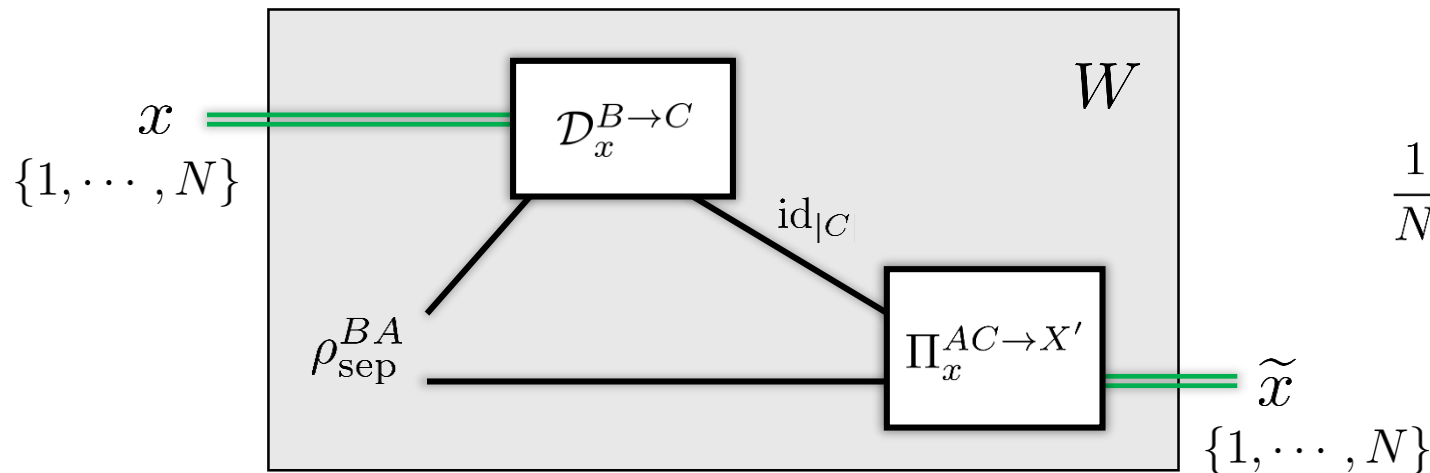
Interpretation: Using any non-entangled state ρ^{AB} and a $|C|$ -dimensional (noiseless) quantum channel, Alice can send Bob N classical messages with $p_{\text{succ}} \leq \min\{\frac{|C|}{N}, 1\}$.



Teleportation and dense coding duality

- The classical threshold of the classical fidelity is $\min\{\frac{|C|}{N}, 1\}$.

Interpretation: Using any non-entangled state ρ^{AB} and a $|C|$ -dimensional (noiseless) quantum channel, Alice can send Bob N classical messages with $p_{\text{succ}} \leq \min\{\frac{|C|}{N}, 1\}$.

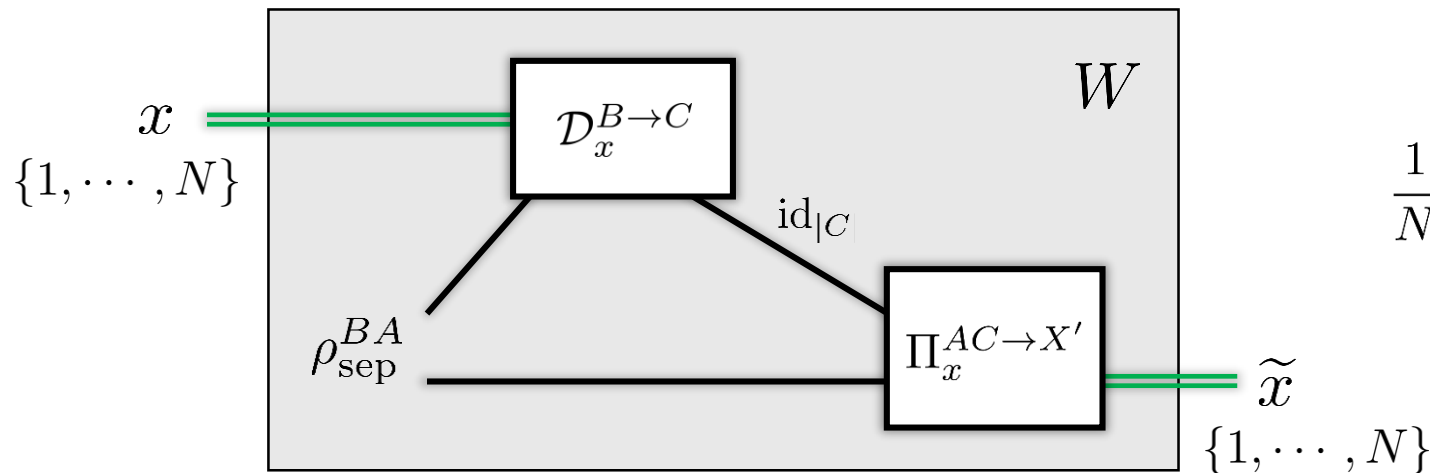


$$\frac{1}{N} \sum_{x=1}^N W(x|x) \leq \min \left\{ \frac{|C|}{N}, 1 \right\}.$$

Teleportation and dense coding duality

- The classical threshold of the classical fidelity is $\min\{\frac{|C|}{N}, 1\}$.

Interpretation: Using any non-entangled state ρ^{AB} and a $|C|$ -dimensional (noiseless) quantum channel, Alice can send Bob N classical messages with $p_{\text{succ}} \leq \min\{\frac{|C|}{N}, 1\}$.



$$\frac{1}{N} \sum_{x=1}^N W(x|x) \leq \min \left\{ \frac{|C|}{N}, 1 \right\}.$$

Chitambar, FL, arXiv:2302.14798

Theorem: For $N > |C|$, a bipartite state ρ^{AB} is useful for dense coding (i.e., can exceed the classical fidelity threshold $|C|/N$) iff there exists a channel $\mathcal{E}^{B \rightarrow C}$ such that $\omega^{AC} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ violates the reduction criterion.

Conclusions and ongoing research

Conclusions and ongoing research

- We have formulated an operational and quantitative duality between teleportation and dense coding, generalizing and extending earlier ideas of Werner.

Werner JPA **34** 7081 (2001)

Conclusions and ongoing research

- We have formulated an operational and quantitative duality between teleportation and dense coding, generalizing and extending earlier ideas of Werner.

Werner JPA **34** 7081 (2001)

- The state discrimination structure of port-based teleportation protocols generalizes to all one-way teleportation protocols.

Conclusions and ongoing research

- We have formulated an operational and quantitative duality between teleportation and dense coding, generalizing and extending earlier ideas of Werner.

Werner JPA **34** 7081 (2001)

- The state discrimination structure of port-based teleportation protocols generalizes to all one-way teleportation protocols.

- One key take-home message:

A bipartite state ρ^{AB} can exceed the classical teleportation threshold
iff it can exceed the classical dense coding threshold.

Conclusions and ongoing research

- We have formulated an operational and quantitative duality between teleportation and dense coding, generalizing and extending earlier ideas of Werner.

Werner JPA **34** 7081 (2001)

- The state discrimination structure of port-based teleportation protocols generalizes to all one-way teleportation protocols.

- One key take-home message:

A bipartite state ρ^{AB} can exceed the classical teleportation threshold
iff it can exceed the classical dense coding threshold.

- What states can exceed these thresholds? How to decide if a given state belongs to this class?

Conclusions and ongoing research

- We have formulated an operational and quantitative duality between teleportation and dense coding, generalizing and extending earlier ideas of Werner. Werner JPA **34** 7081 (2001)

- The state discrimination structure of port-based teleportation protocols generalizes to all one-way teleportation protocols.

- One key take-home message:

A bipartite state ρ^{AB} can exceed the classical teleportation threshold
iff it can exceed the classical dense coding threshold.


- What states can exceed these thresholds? How to decide if a given state belongs to this class?
- We need to decide if there exists a channel $\mathcal{E}^{B \rightarrow C}$ such that $\omega^{AC} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ violates the reduction criterion.

Conclusions and ongoing research

Conclusions and ongoing research

- We can phrase this question as a bilinear optimization problem.

Conclusions and ongoing research

- We can phrase this question as a bilinear optimization problem.
- For a given state ρ^{AB} , define the conditional state $\rho^{B|A} = (\rho^A)^{-1/2} \rho^{AB} (\rho^A)^{-1/2}$:
 $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion iff $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B}) \leq \mathbb{I}^A \otimes \mathbb{I}^B$.

Conclusions and ongoing research

- We can phrase this question as a bilinear optimization problem.

- For a given state ρ^{AB} , define the conditional state $\rho^{B|A} = (\rho^A)^{-1/2} \rho^{AB} (\rho^A)^{-1/2}$:

 $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion iff $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B}) \leq \mathbb{I}^A \otimes \mathbb{I}^B$.

- The largest eigenvalue of $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B})$ is $\max_{|\varphi\rangle^{AC}} \langle \varphi | \text{Tr}_B((\rho^{B|A})^\top J_{\mathcal{E}}) | \varphi \rangle$,

where $J_{\mathcal{E}} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\phi^{+B\tilde{B}})$ is the Choi matrix of \mathcal{E} .

Conclusions and ongoing research

- We can phrase this question as a bilinear optimization problem.

- For a given state ρ^{AB} , define the conditional state $\rho^{B|A} = (\rho^A)^{-1/2} \rho^{AB} (\rho^A)^{-1/2}$:

⇒ $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion iff $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B}) \leq \mathbb{I}^A \otimes \mathbb{I}^B$.

- The largest eigenvalue of $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B})$ is $\max_{|\varphi\rangle^{AC}} \langle \varphi | \text{Tr}_B((\rho^{B|A})^\top J_{\mathcal{E}}) | \varphi \rangle$,

where $J_{\mathcal{E}} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\phi^{+B\tilde{B}})$ is the Choi matrix of \mathcal{E} .

- Then ρ^{AB} can exceed the classical teleportation threshold iff

$$1 < \max_{J^{BC}} \max_{\omega^{\tilde{A}C}} \text{Tr}[\omega^{\tilde{A}C} \text{Tr}_B((\rho^{B|A})^\top J^{BC})]$$

$$\text{subject to } \text{Tr}_C J^{BC} = \mathbb{I}^B;$$

$$\text{Tr}[\omega] = 1;$$

$$\omega, J \geq 0.$$

Conclusions and ongoing research

- We can phrase this question as a bilinear optimization problem.
- For a given state ρ^{AB} , define the conditional state $\rho^{B|A} = (\rho^A)^{-1/2} \rho^{AB} (\rho^A)^{-1/2}$:

⇒ $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion iff $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B}) \leq \mathbb{I}^A \otimes \mathbb{I}^B$.

- The largest eigenvalue of $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B})$ is $\max_{|\varphi\rangle^{AC}} \langle \varphi | \text{Tr}_B((\rho^{B|A})^\top J_{\mathcal{E}}) | \varphi \rangle$,

where $J_{\mathcal{E}} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\phi^{+B\tilde{B}})$ is the Choi matrix of \mathcal{E} .

- Then ρ^{AB} can exceed the classical teleportation threshold iff

$$1 < \max_{J^{BC}} \max_{\omega^{\tilde{A}C}} \text{Tr}[\omega^{\tilde{A}C} \text{Tr}_B((\rho^{B|A})^\top J^{BC})]$$

$$\text{subject to } \text{Tr}_C J^{BC} = \mathbb{I}^B;$$

$$\text{Tr}[\omega] = 1;$$

$$\omega, J \geq 0.$$

– Numerically feasible for small dimensions

Conclusions and ongoing research

- We can phrase this question as a bilinear optimization problem.
- For a given state ρ^{AB} , define the conditional state $\rho^{B|A} = (\rho^A)^{-1/2} \rho^{AB} (\rho^A)^{-1/2}$:

⇒ $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion iff $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B}) \leq \mathbb{I}^A \otimes \mathbb{I}^B$.

- The largest eigenvalue of $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B})$ is $\max_{|\varphi\rangle^{AC}} \langle \varphi | \text{Tr}_B((\rho^{B|A})^\top J_{\mathcal{E}}) | \varphi \rangle$,

where $J_{\mathcal{E}} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\phi^{+B\tilde{B}})$ is the Choi matrix of \mathcal{E} .

- Then ρ^{AB} can exceed the classical teleportation threshold iff

$$1 < \max_{J^{BC}} \max_{\omega^{AC}} \text{Tr}[\omega^{\tilde{A}C} \text{Tr}_B((\rho^{B|A})^\top J^{BC})]$$

$$\text{subject to } \text{Tr}_C J^{BC} = \mathbb{I}^B;$$

$$\text{Tr}[\omega] = 1;$$

$$\omega, J \geq 0.$$

– Numerically feasible for small dimensions

– Current work involves analyzing this for certain families of states.

Conclusions and ongoing research

- We can phrase this question as a bilinear optimization problem.
- For a given state ρ^{AB} , define the conditional state $\rho^{B|A} = (\rho^A)^{-1/2} \rho^{AB} (\rho^A)^{-1/2}$:

 $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{AB})$ satisfies the reduction criterion iff $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B}) \leq \mathbb{I}^A \otimes \mathbb{I}^B$.

- The largest eigenvalue of $\text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\rho^{A|B})$ is $\max_{|\varphi\rangle^{AC}} \langle \varphi | \text{Tr}_B((\rho^{B|A})^\top J_{\mathcal{E}}) | \varphi \rangle$,

where $J_{\mathcal{E}} = \text{id}^A \otimes \mathcal{E}^{B \rightarrow C}(\phi^{+B\tilde{B}})$ is the Choi matrix of \mathcal{E} .

- Then ρ^{AB} can exceed the classical teleportation threshold iff

$$1 < \max_{J^{BC}} \max_{\omega^{AC}} \text{Tr}[\omega^{\tilde{A}C} \text{Tr}_B((\rho^{B|A})^\top J^{BC})]$$

$$\text{subject to } \text{Tr}_C J^{BC} = \mathbb{I}^B;$$

$$\text{Tr}[\omega] = 1;$$

$$\omega, J \geq 0.$$

– Numerically feasible for small dimensions

– Current work involves analyzing this for certain families of states.

– Can it be reduced to an SDP?

Thank you for your attention!

